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# Physics of transient power loads on plasma facing components in ITER

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ITER Organization

# Outline

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- Introduction
  - ✓ Plasma regimes in ITER
  - ✓ Introduction to major phenomena leading to transient power fluxes in ITER and consequences for plasma facing components (PFCs)
- Power fluxes to PFCs during ELMs and methods for ELM control/suppression
- Power fluxes to PFCs during disruptions and disruption mitigation schemes
- Summary and Conclusions

# Plasmas in ITER (I)

ITER Mission : “To demonstrate the scientific and technological feasibility of fusion energy for peaceful purposes”

➤ ITER major fusion performance goal

$$\alpha \text{ dominated plasmas } (P_{\alpha}/P_{\text{add}} = 2 \leftrightarrow Q_{\text{DT}} = 10)$$

&

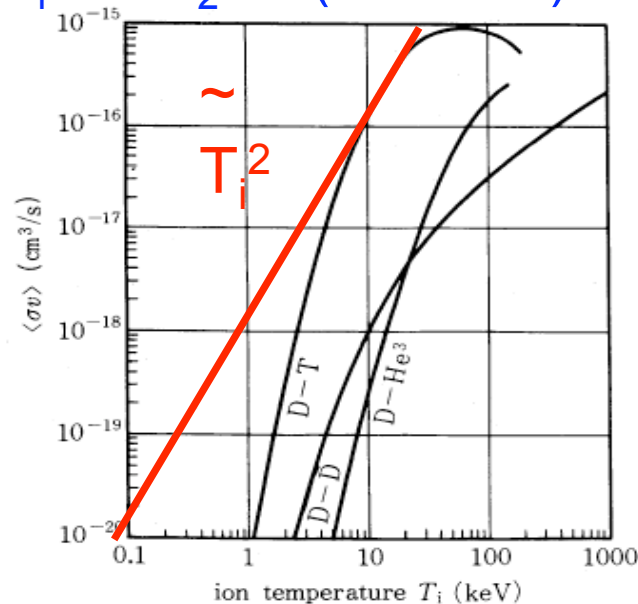
$$P_{\text{fusion}} = 500 \text{ MW}$$

- ✓ Inductive operation with 300-500 s burn time
- ✓ H-mode energy/particle confinement H-mode

## Plasmas in ITER (II)

ITER transient power fluxes  $\leftrightarrow$  ITER fusion goal

- $P_{\text{fusion}} \leftrightarrow$  Plasma thermal energy :  $W_{\text{plasma}} = \frac{3}{2} \langle n_e T_e + n_i T_i \rangle V_{\text{plasma}}$   
 ${}^2_1\text{D} + {}^3_1\text{T} \rightarrow {}^4_2\text{He} (3.5 \text{ MeV}) + {}^1_0\text{n} (14.1 \text{ MeV})$



$$P_{\text{fusion}} = E_{\text{fusion}} \langle n_D \rangle \langle n_T \rangle \langle \sigma v \rangle_{DT} V_{\text{plasma}} \sim \langle n_{DT} \rangle^2 \langle T_{DT} \rangle^2 V_{\text{plasma}} \sim W_{\text{plasma}}^2$$

$$P_{\text{fusion}} = 500 \text{ MW} \leftrightarrow W_{\text{plasma}} = 350 \text{ MJ}$$

$$W_{\text{magnetic}} = \frac{1}{2} L_{\text{plasma}} I_p^2 \quad \text{with} \quad L_{\text{plasma}} = \mu_0 R_0 \left( \ln \left( 8 \frac{R_0}{a} \right) - 2 + \frac{\ell_i}{4} + \frac{3}{4} \beta_p \right) \quad \sim 5 W_{\text{plasma}}$$

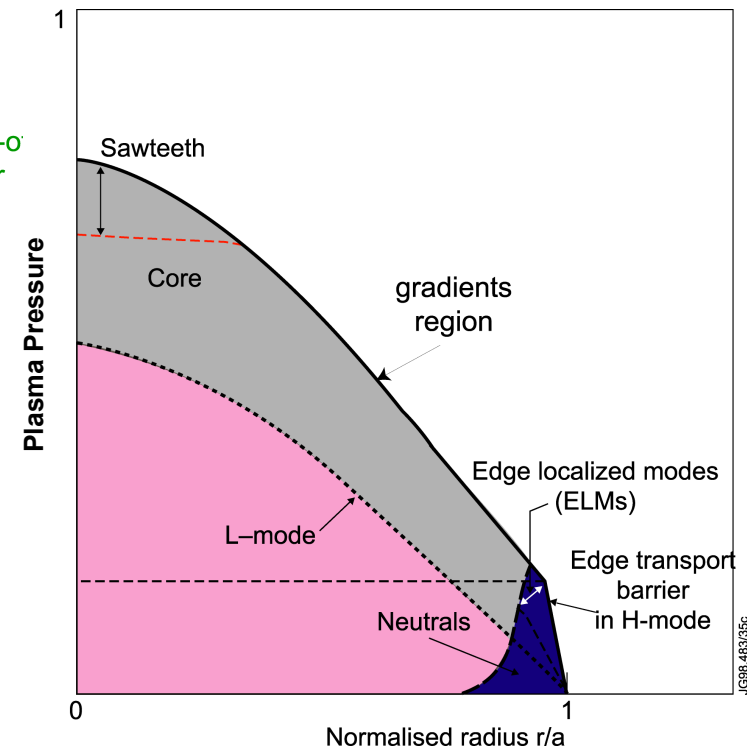
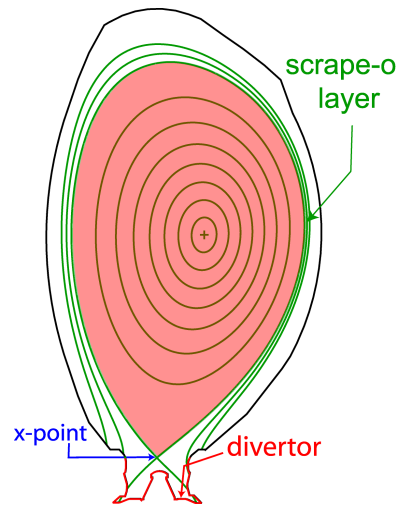
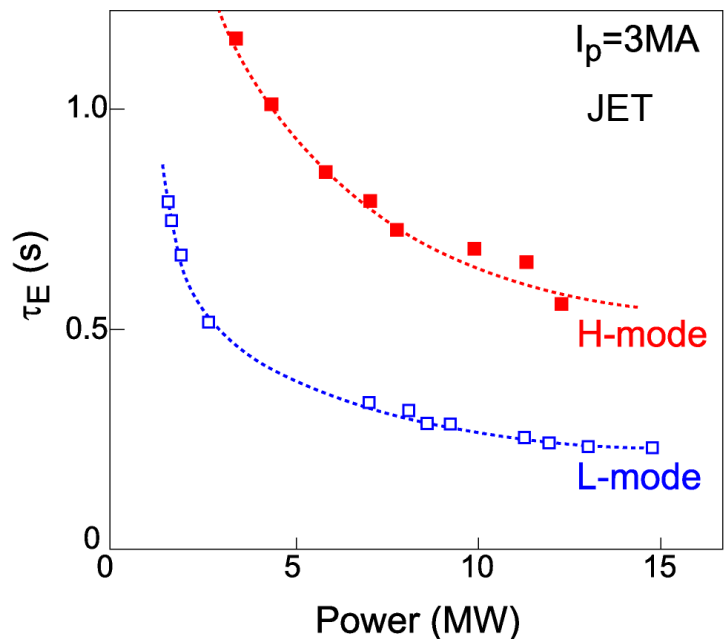
# Plasmas in ITER (III)

## ITER transient power fluxes $\leftrightarrow$ confinement regime

Energy Confinement in Tokamaks and Stellarators : **L** (*low*) and **H** (*high*)  
Confinement Modes

H-mode  $\leftrightarrow$  **E**dge **T**ransport **B**arrier ( $\rightarrow$  Pedestal)

In Tokamaks :  $P_{\text{INPUT}} > P_{\text{L-H}}(n_e, B_t, R)$



$$\tau_E = W_{\text{plasma}} / P_{\text{input}} \rightarrow W_{\text{plasma}} \uparrow^2 \ \& \ P_{\text{fusion}} \sim W_{\text{plasma}}^2 \uparrow^4$$

# Phenomena causing transients in ITER (I)

➤ Largest energy transients in ITER → disruptions

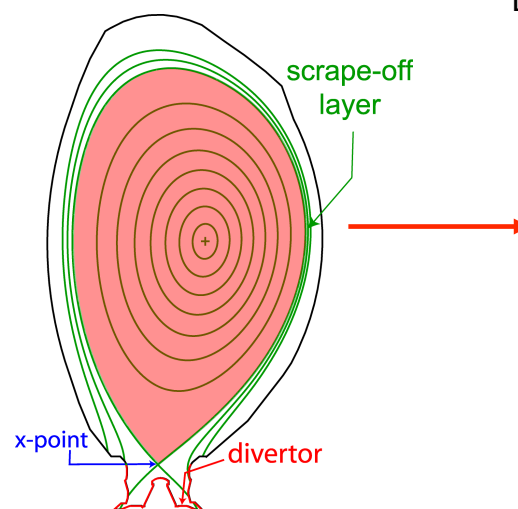
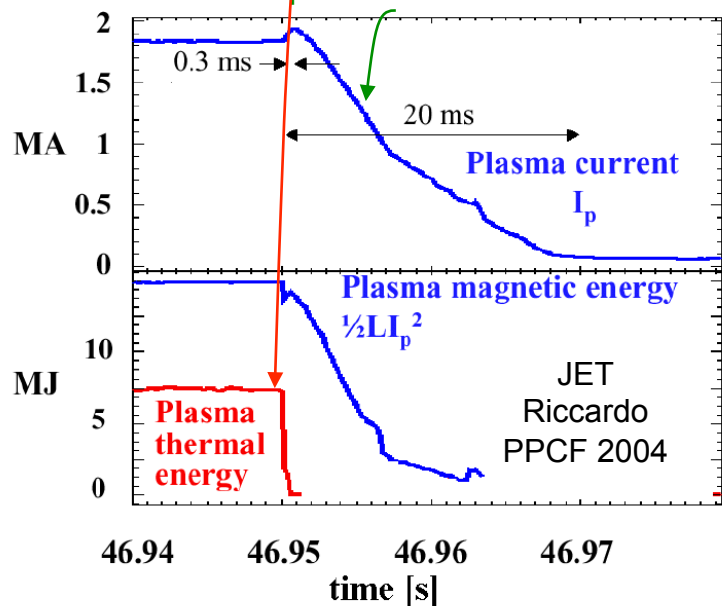
✓  $W_{\text{plasma}}$  → deposited by plasma onto PFCs

✓  $W_{\text{magnetic}}$  → conductors & VV + radiation/plasma onto PFCs or high  $E_e$

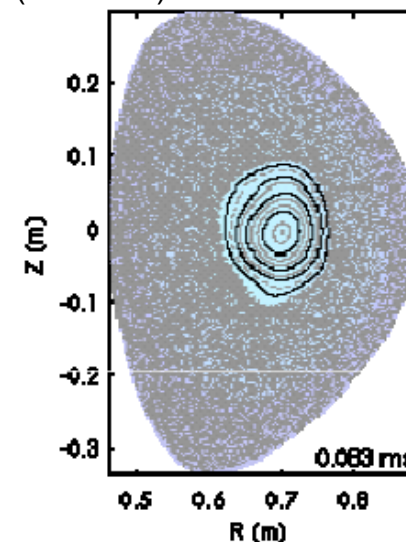
Plasma develops unstable  $p(r)$ ,  $j(r)$  → Large scale MHD unstable modes grow

plasma confinement is destroyed (thermal quench ~ ms) →  $W_{\text{plasma}}$

plasma current vanishes (current quench ~10s ms) →  $W_{\text{magnetic}}$



DIID-D (NIMROD) – V. Izzo NF 2007

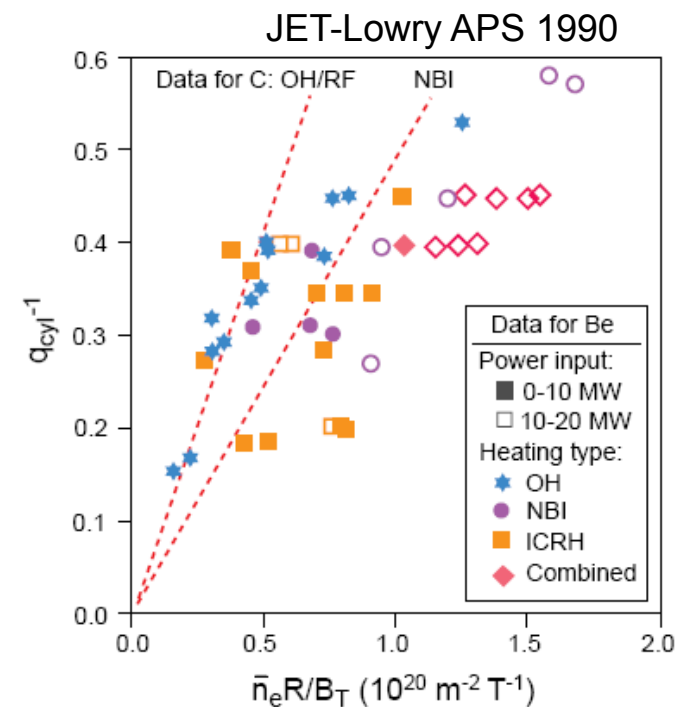
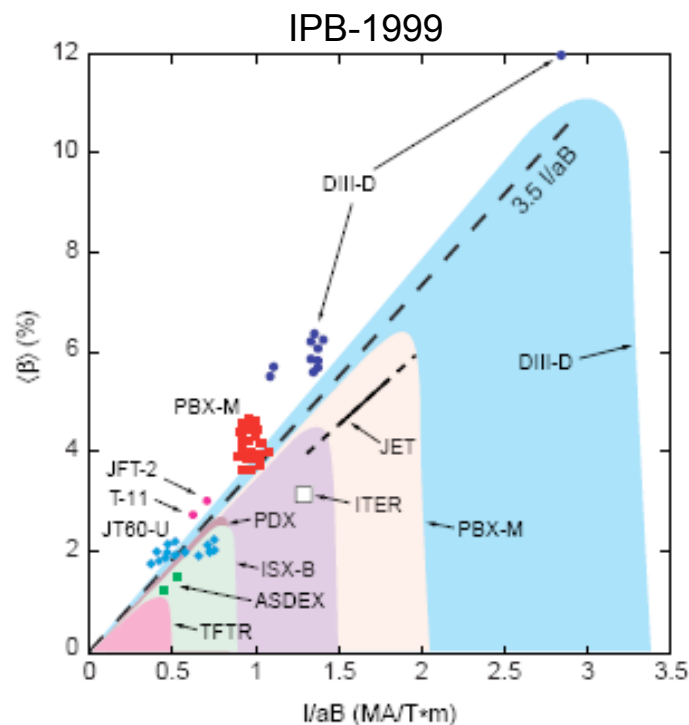


# Phenomena causing transients in ITER (II)

Tokamak operation is typically disruption-limited by three limits :

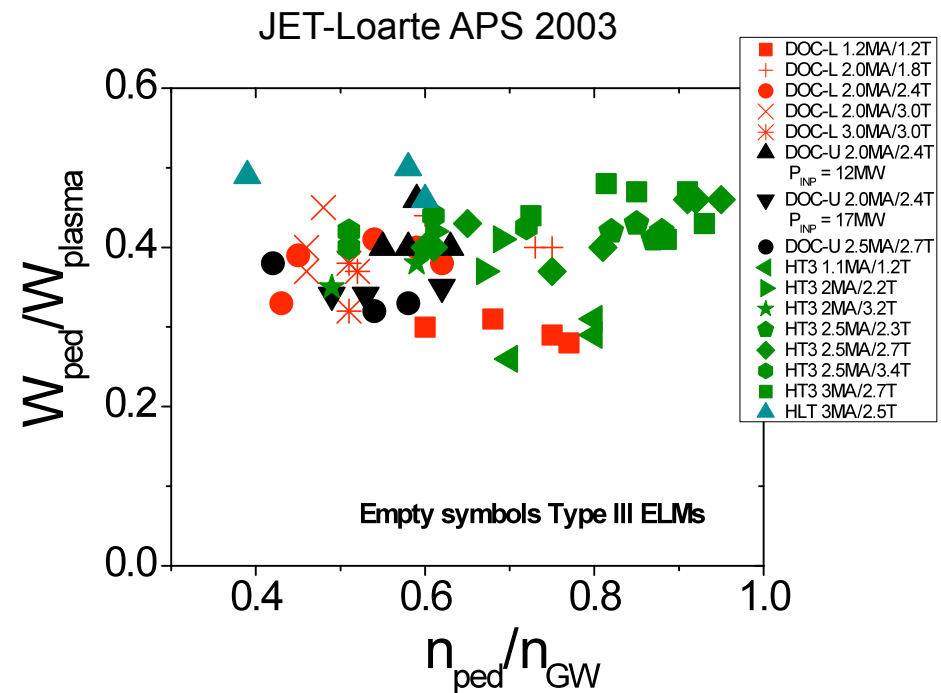
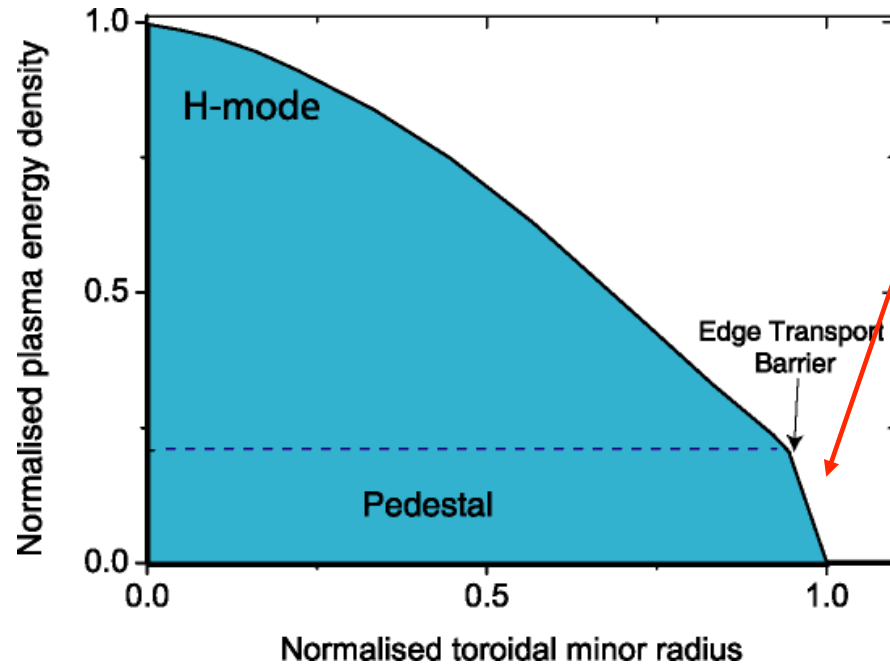
- Pressure limit  $\rightarrow P_{\text{fusion}} \sim \beta^2$
- Density + Radiation limit  $\rightarrow T < 30\text{-}40 \text{ keV}$  & stationary power flux control
- Low q limit  $\rightarrow q \sim B_t/I_p$

Operation of ITER (and tokamak fusion reactors) approaches these limits



# Phenomena causing transients in ITER (III)

High  $W_{\text{plasma}}$  in H-mode associated with pedestal  $\rightarrow$  large grad-p and  $j_{\text{edge}}$



In H-mode :  $W_{\text{plasma}} \sim W_{\text{ped}}$

ITER :  $W_{\text{ped}} \sim 100 - 140 \text{ MJ}$

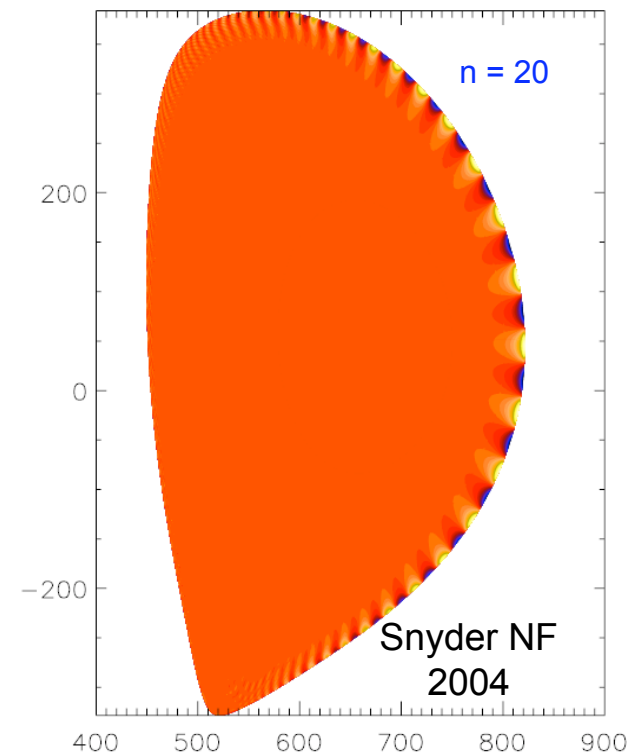
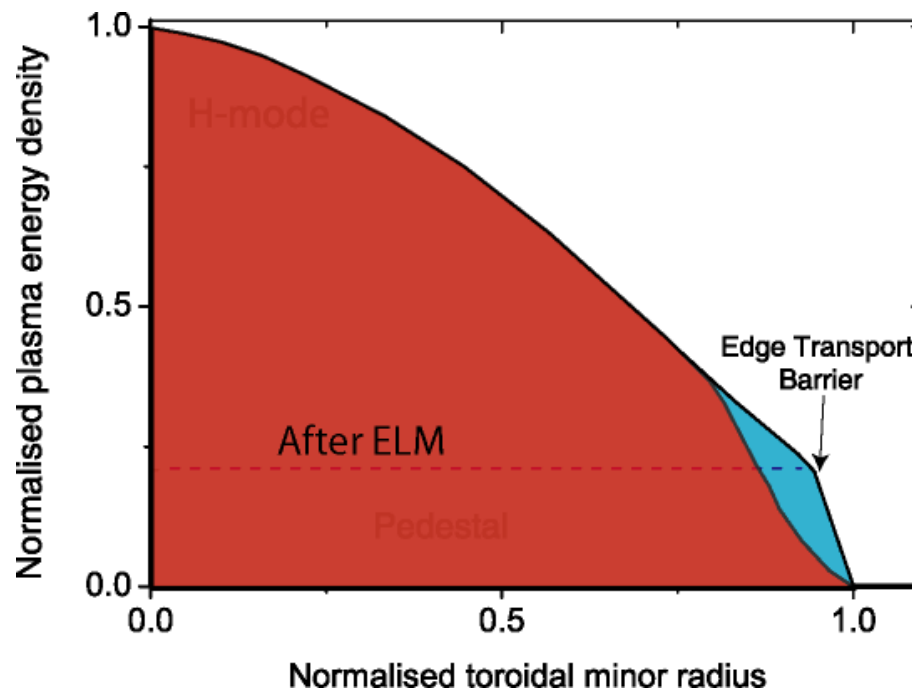
$n_{\text{ped}} = 7-9 \cdot 10^{19} \text{ m}^{-3}$ ,  $T_{\text{ped}} = 3-5 \text{ keV}$

$$W_{\text{ped}} = \frac{3}{2} n_{\text{ped}} (T_{e,\text{ped}} + T_{i,\text{ped}}) V_{\text{plasma}}$$



# Phenomena causing transients in ITER (IV)

- Edge Transport Barrier & ELMs : large edge grad-p(r) & j(r) → edge MHD instability → quasi-periodic relaxations (ELMs) →  $\Delta W_{\text{ELM}}$
- Various Types of ELMs (grad-p,  $j_{\text{edge}}$ ): highest  $W_{\text{plasma}}$  vs.  $n_e$  and “operating space” Type I ELMs → basis for ITER  $Q_{\text{DT}} = P_{\text{fus}}/P_{\text{inp}} = 10$



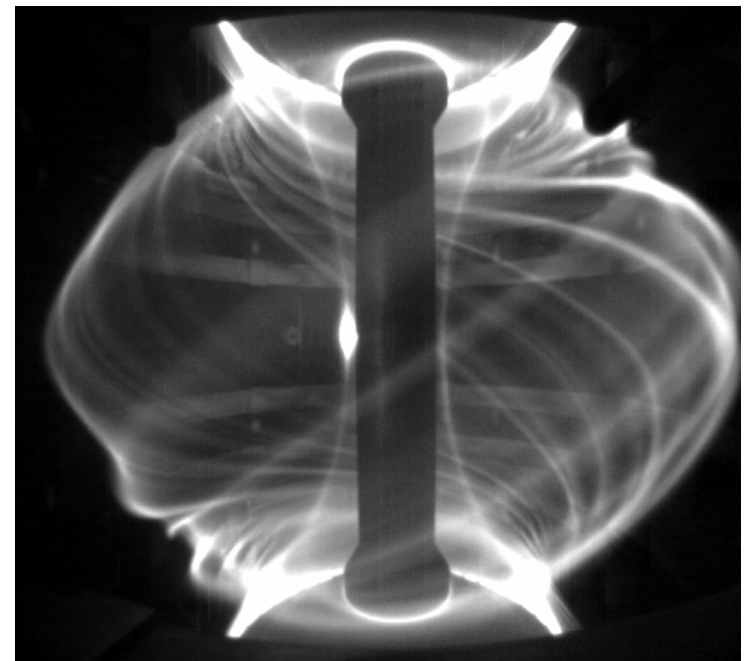
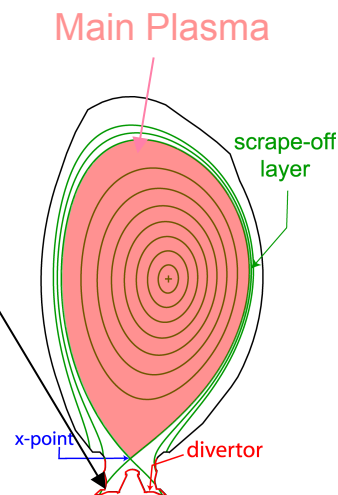
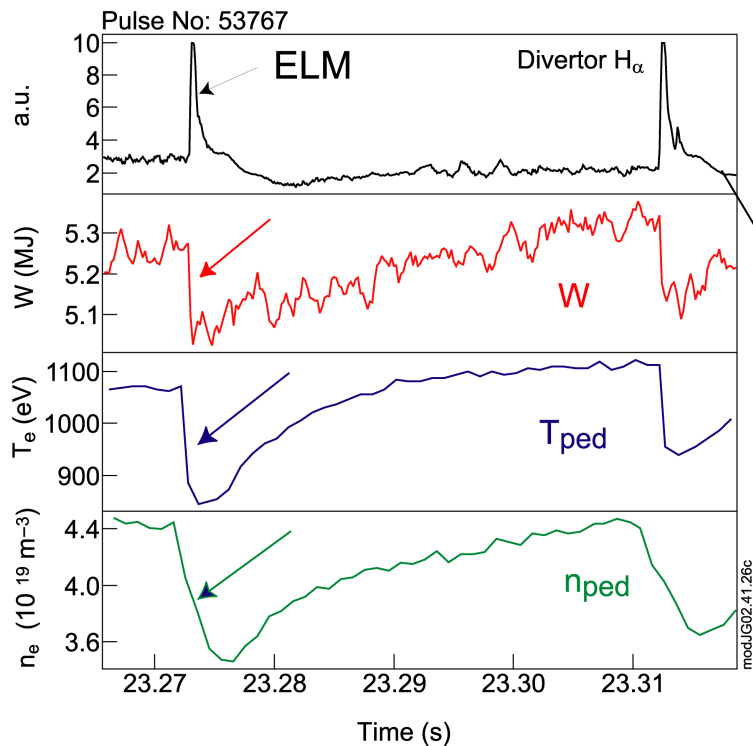
# Phenomena causing transients in ITER (V)

Exterior region of plasma ( $r/a < 0.75$ ) experiences quasi-periodic relaxations during ELMs

$\Delta W_{ELM}$  small Fraction of  $W_{plasma}$  ( $< 10\%$ ) in  $\sim 200 \mu s \rightarrow$  Large Energy Flux

JET – Type I ELM - Loarte PPCF 2002

MAST- Kirk, EPS'06

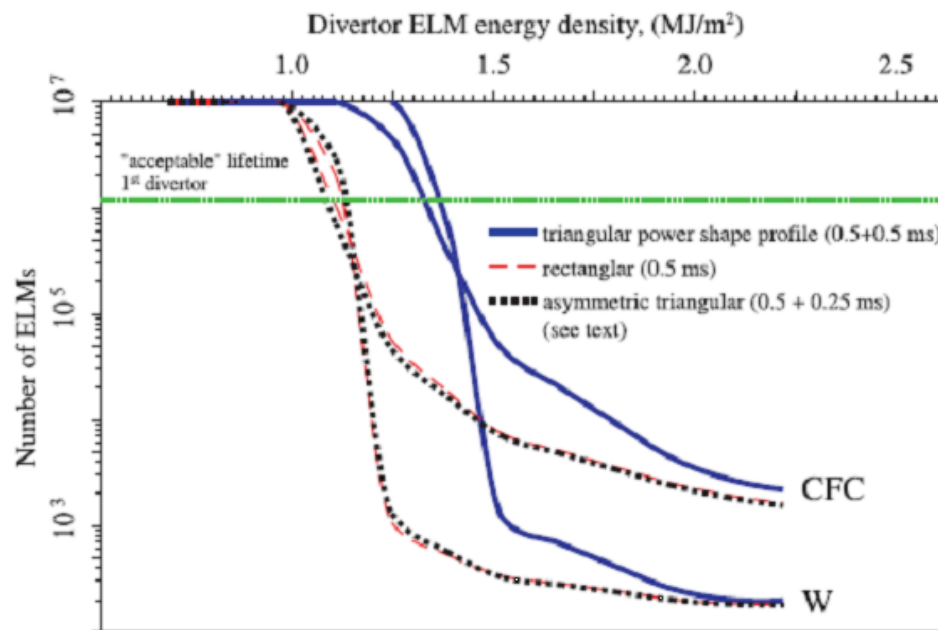


# Consequences for ITER PFCs (I)

- Magnitude of transient power fluxes in ITER → Deterioration of PFCs by large erosion and deformation (W melting)
  - ✓ Reduced lifetime PFC lifetime
  - ✓ Difficulties to operation on damaged PFCs (hot spots)
  - ✓ Generation of dust, ...
- For events lasting  $\sim 1$  ms → damage threshold  $< 1 \text{ MJm}^{-2}$  (CFC & W)

$\sim 100 Q_{DT} = 10$   
discharge

$\sim 1 Q_{DT} = 10$   
discharge

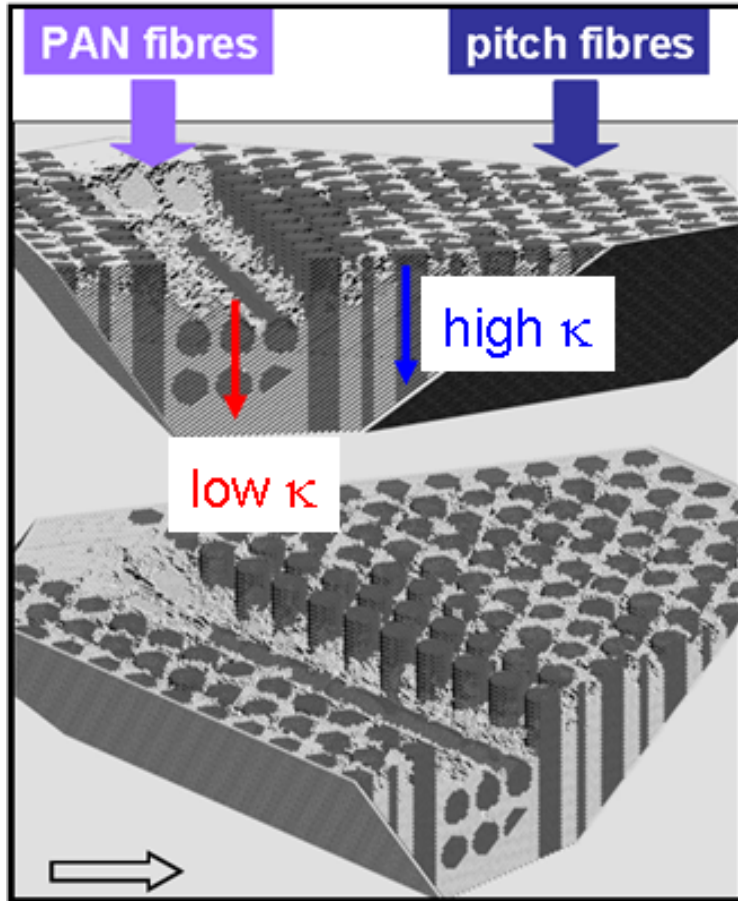


Federici PPCF 2003

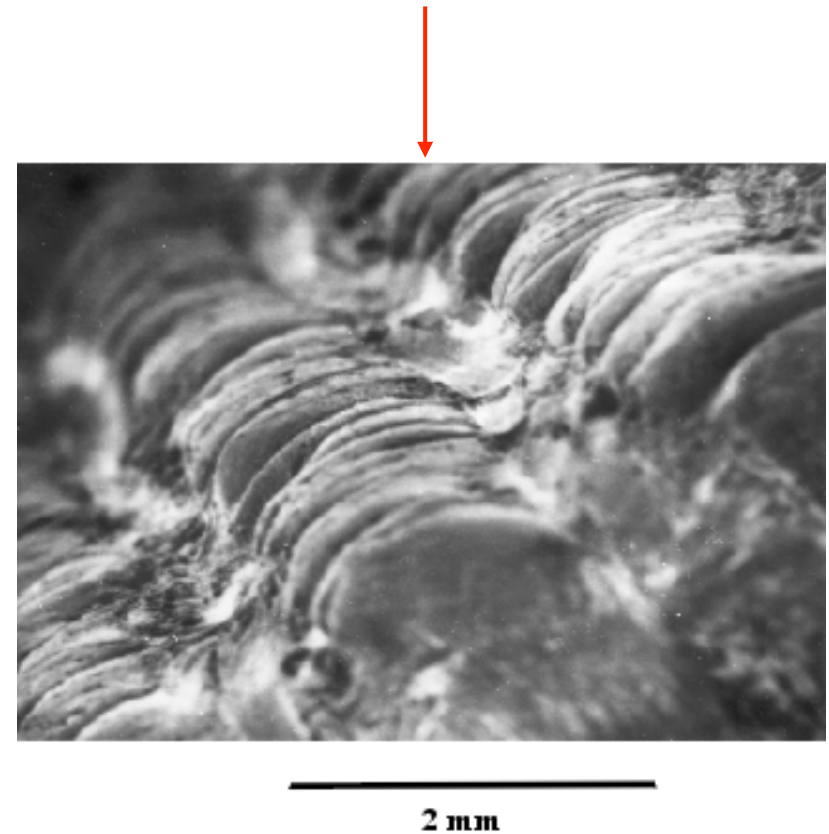
# Consequences for ITER PFCs (II)

## ➤ Experimental studies of material damage under ITER-like transients

S. Pestchanyi JNM 2007

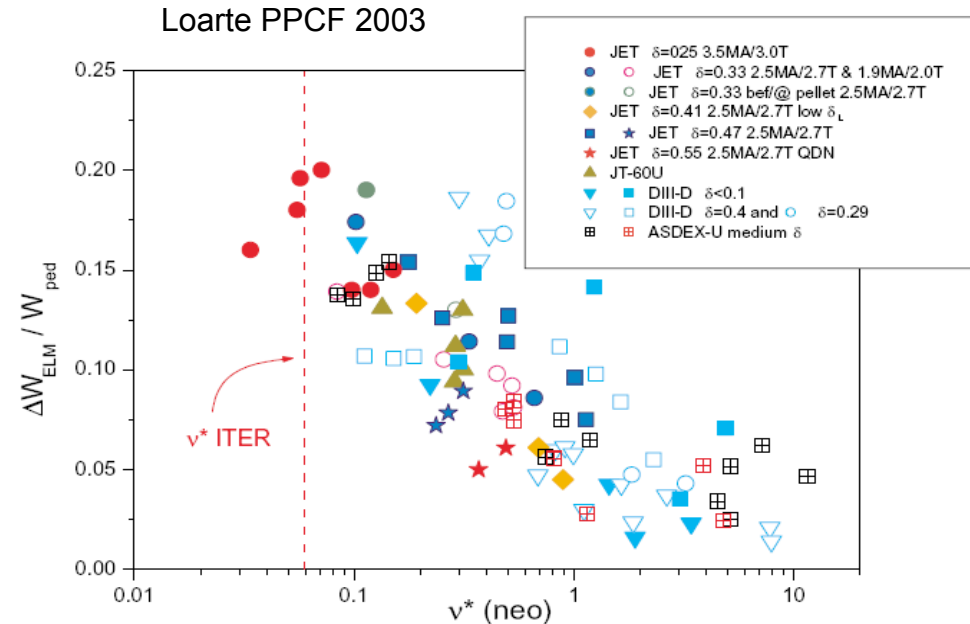
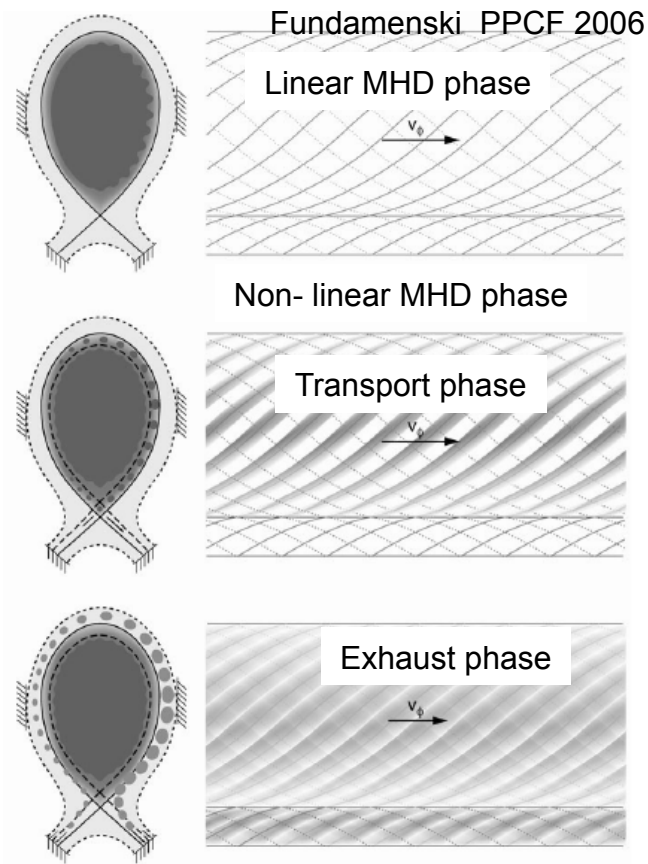


CFC target exposed to ITER-like loads in Plasma-Gun experiments



# Power fluxes to PFCs during ELMs (I)

- Basic physics picture of processes leading to energy loss by ELMs ( $\Delta W_{ELM}$ ) well developed
- Quantitative modelling/prediction of  $\Delta W_{ELM}$  outstanding  $\rightarrow$  empirical extrapolation to ITER

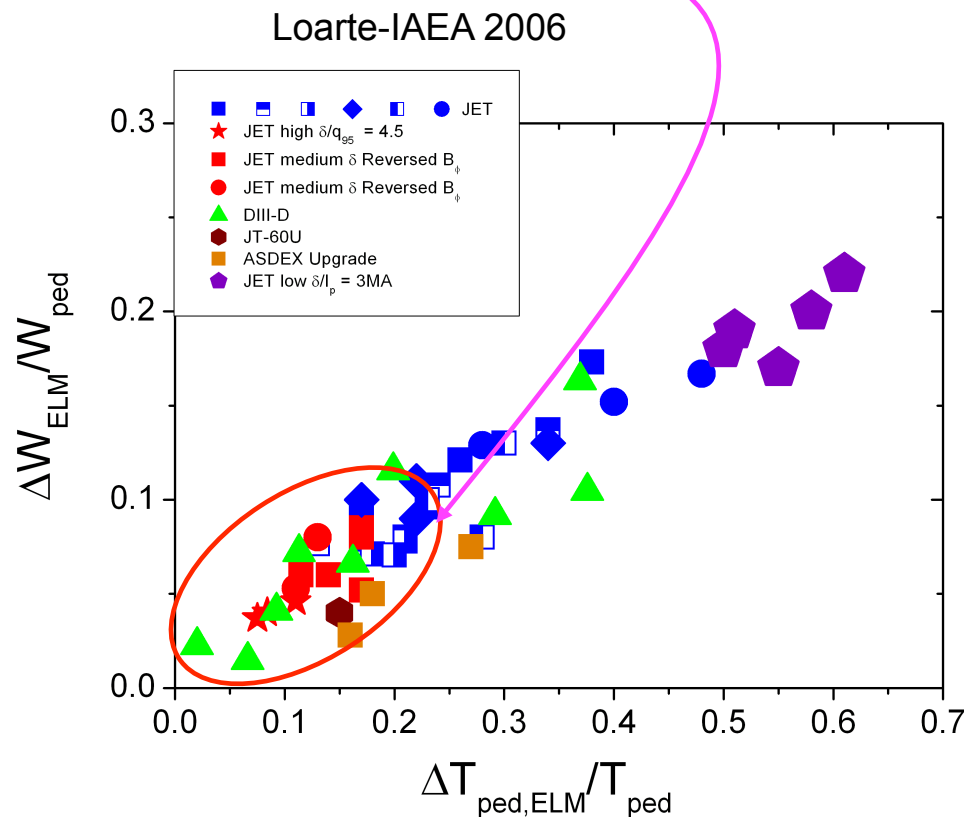
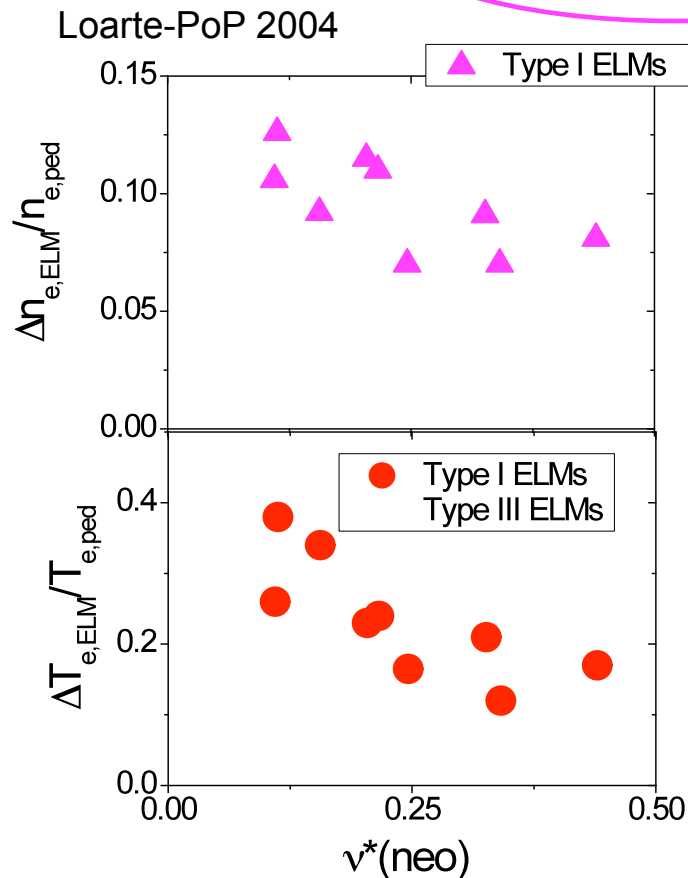


$$v^* \sim Rn_e / T_e^2$$

ITER :  $W_{ped} \sim 100-140$  MJ  $\rightarrow \Delta W_{ELM} > 20$  MJ

## Power fluxes to PFCs during ELMs (II)

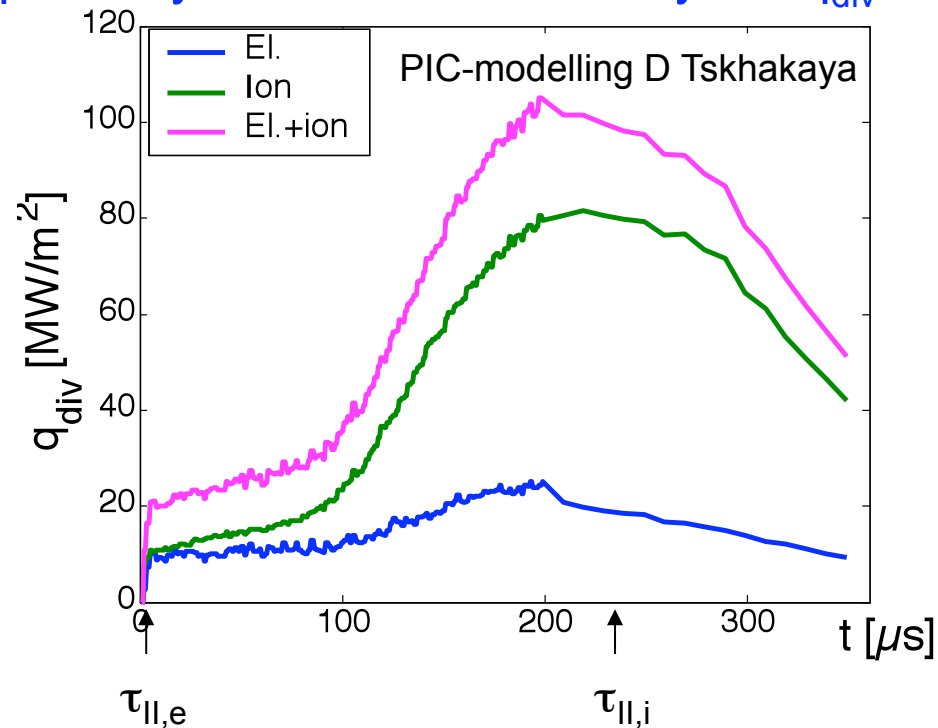
$$\Delta W_{ELM} = (3 \langle n_{ped} \rangle \Delta T_{ped,ELM} + 3 \langle T_{ped} \rangle \Delta n_{ped,ELM}) V_{ELM}$$



➤ Reduction of  $\Delta W_{ELM}/W_{ped}$  mainly by reduction of  $\Delta T_{ELM}/T_{ped}$  → decrease of conductive ELM losses → R&D decrease of conductive losses compatible with ITER requirements

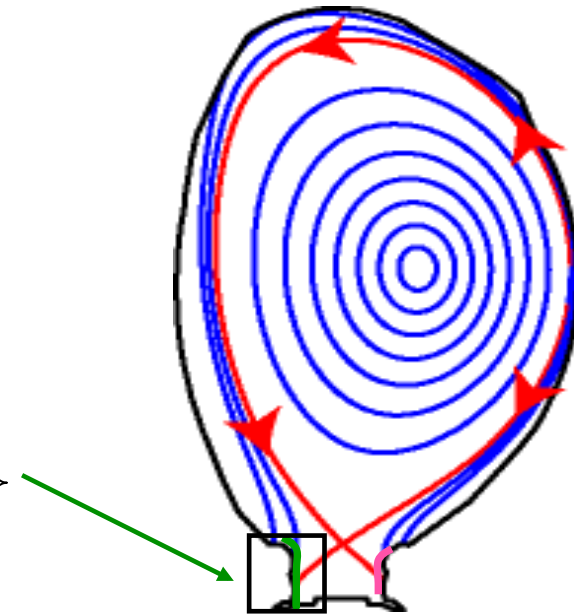
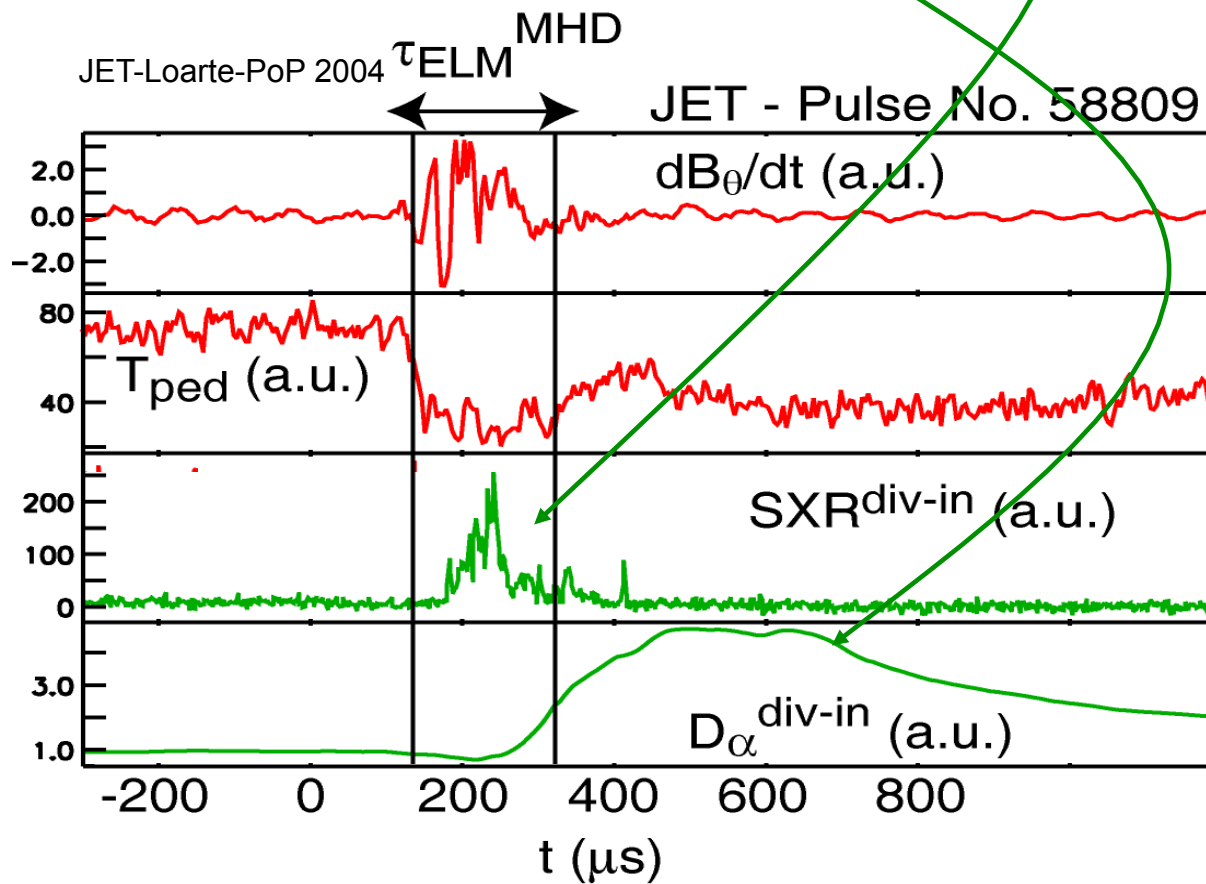
## Power fluxes to PFCs during ELMs (III)

- In non-linear MHD phase plasma with  $n \sim n_{ped}$ ,  $T \sim T_{ped}$  connects to PFC along open field lines
  - ✓ Hot electrons arrive at the divertor target  $v_e/v_i \sim (m_i/m_e)^{1/2} \sim 60$  ( $\tau \sim \mu s$ )
  - ✓ Formation of sheath with  $T \sim T_{ped} \rightarrow$  acceleration of ions
  - ✓ Ions arrive at target in timescale of  $\tau_{||}$  ( $\sim 100$ 's of  $\mu s$ )
  - ✓ ELM power pulse dynamics dominated by ion  $q_{div}^{ELM}(t) \sim q_{ion}^{ELM}(t)$



# Power fluxes to PFCs during ELMs (IV)

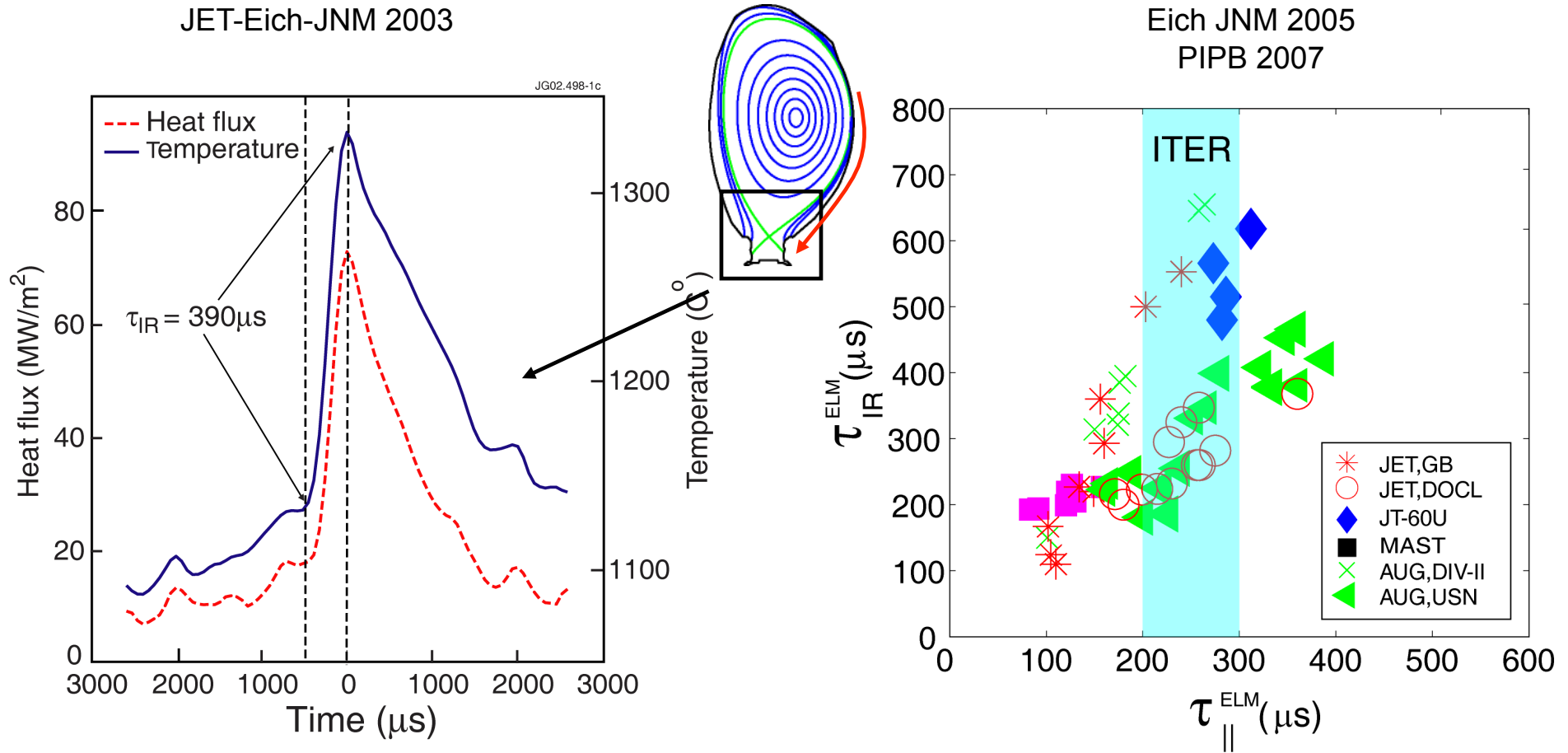
- Basic ELM power flux physics picture experimental confirmation
  - ✓ Hot electrons soft X-ray emission
  - ✓ Delayed ion arrival ( $D_\alpha$ )





# Power fluxes to PFCs during ELMs (V)

- Time scale of divertor ELM energy flux rise correlated with  $\tau_{||,i} \sim L/v_i(T_{ped})$

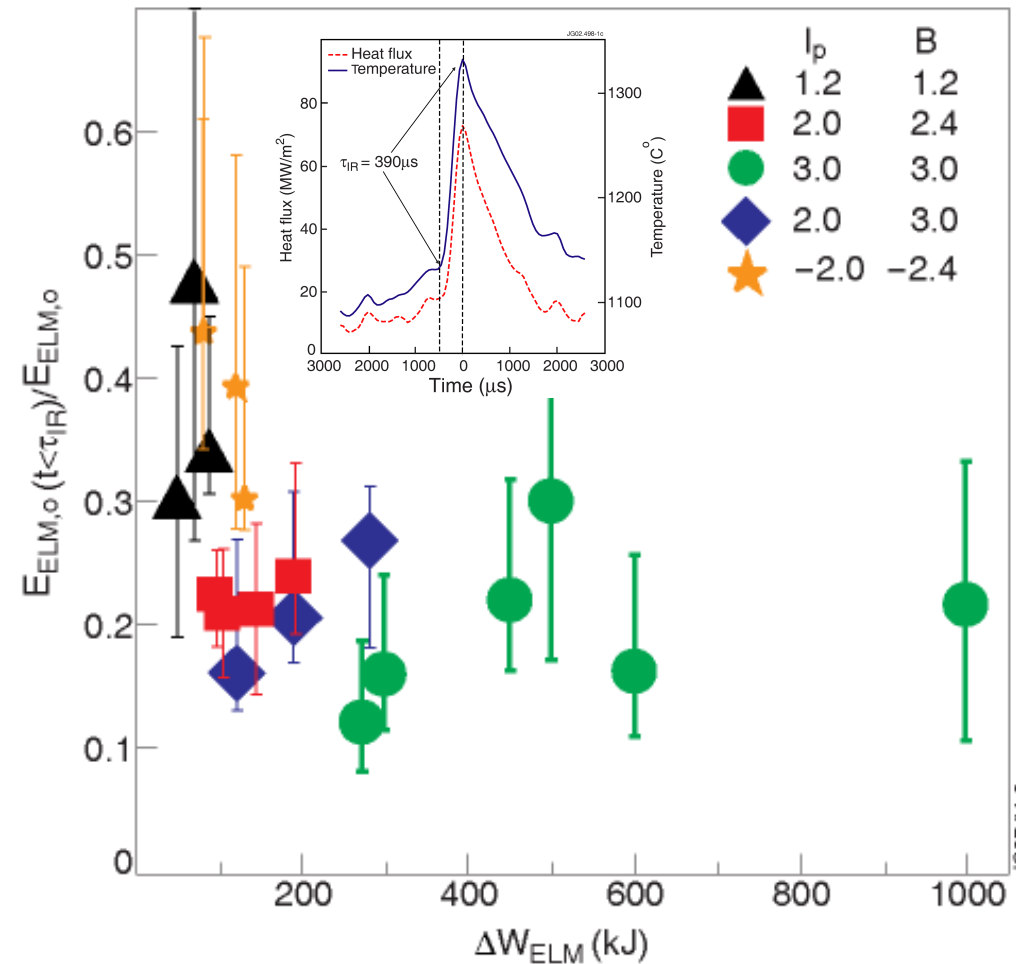


Physics basis for ELM power flux duration in ITER through plasma conditions  $(n_{ped}, T_{ped})$  &  $R \rightarrow R\&D \tau_{ELM}^{IR} \sim \tau_{||}$  relation (pre-ELM divertor plasma, etc.)

# Power fluxes to PFCs during ELMs (VI)

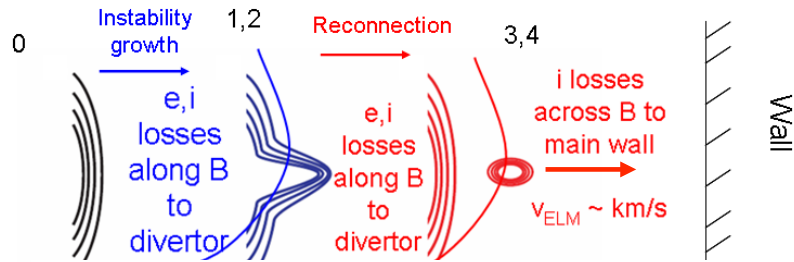
- Because of sheath formation → large proportion of  $\Delta W_{ELM}$  arrives after  $\tau_{IR}$  → smaller  $\Delta T_{surf}$  for given  $\Delta W_{ELM}$

JET-Eich-JNM 2005, Pitts-IAEA 2006

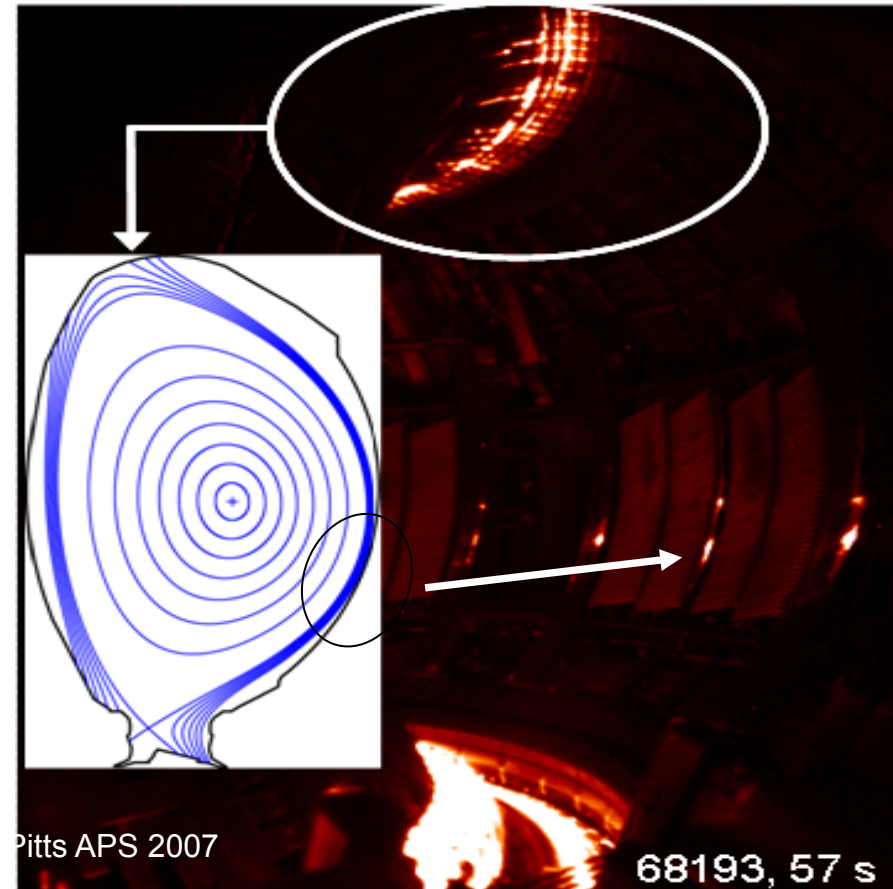
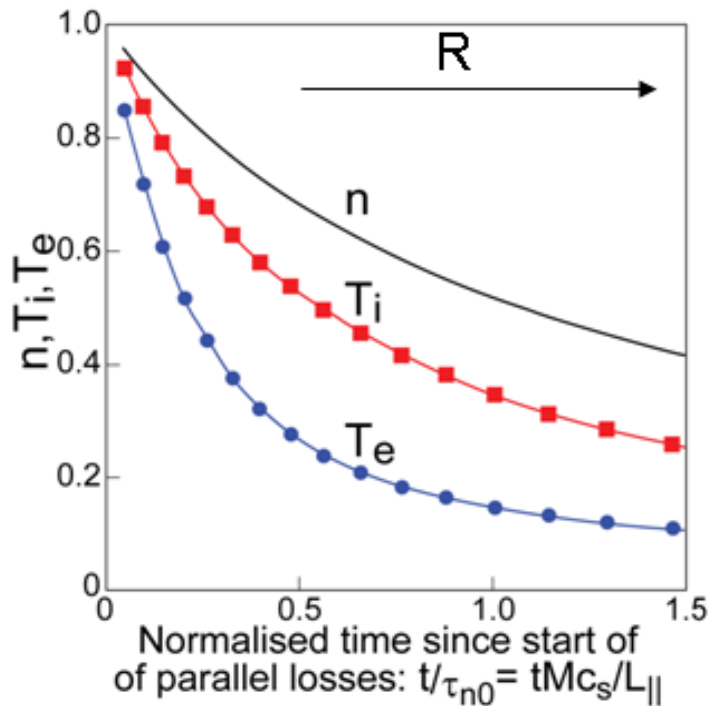


# Power fluxes to PFCs during ELMs (VII)

- ELM energy transport and MHD origin determine areas for deposition of power and power sharing between PFCs : e + i near separatrix & i



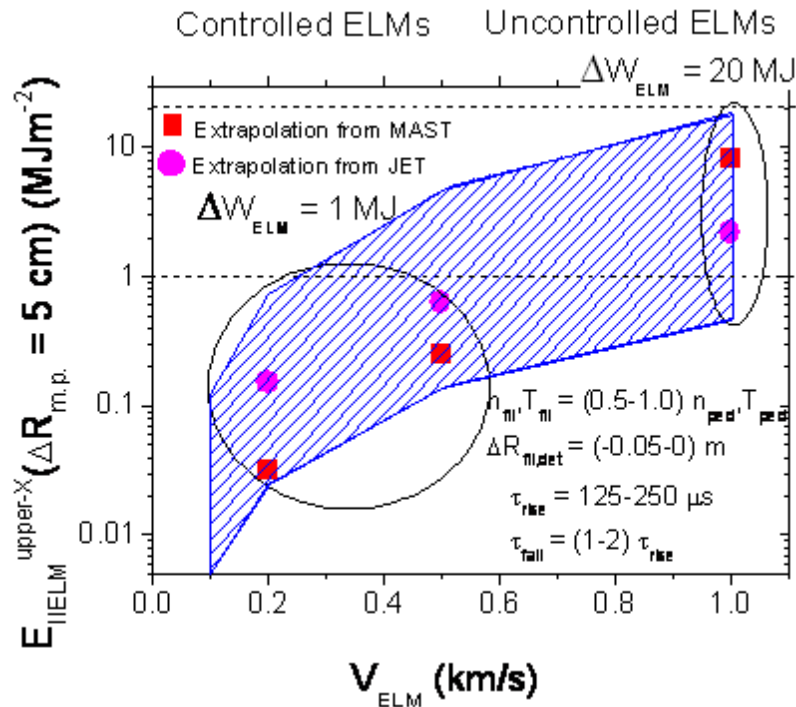
Fundamenski - PPCF 2006



## Power fluxes to PFCs during ELMs (VIII)

- Physics model and physics based extrapolations applied to determine ELM power fluxes to divertor and wall in ITER (R&D on-going)

| ELM Type     | $\Delta W_{\text{ELM}}$<br>(MJ) | $E_{\parallel, \text{max}}^{\text{div-in}}$<br>(MJm <sup>-2</sup> ) | $E_{\text{max}}^{\text{div-in}}$<br>(MJm <sup>-2</sup> ) | $E_{\parallel, \text{max}}^{\text{div-out}}$<br>(MJm <sup>-2</sup> ) | $E_{\text{max}}^{\text{div-out}}$<br>(MJm <sup>-2</sup> ) | $\tau_{\text{ELM}}^{\text{rise}}$<br>( $\mu\text{s}$ ) |
|--------------|---------------------------------|---|--|--|---|--|
| uncontrolled | 20                              | 300   | 17   | 180  | 8.5   | 250-500  |



- Uncontrolled ELMs in ITER → energy fluxes exceed by large factors ( $\sim 10$ ) material damage thresholds → damage to PFC and lifetime reduction
- Control of ELM power fluxes is mandatory for reliable ITER operation

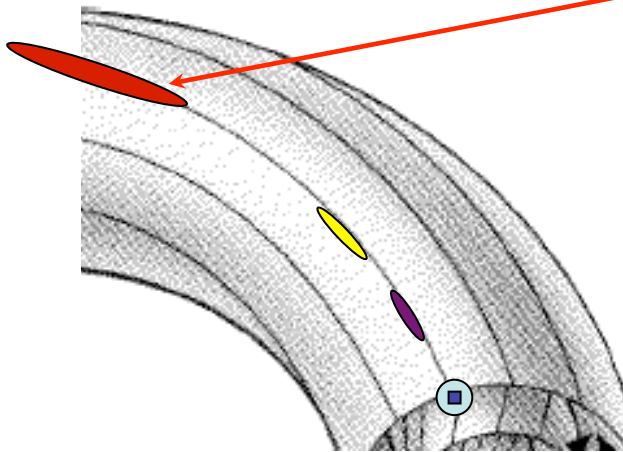
# Active control of ELMs (I)

Two methods considered for ELM control in ITER → injection of frozen pellets & perturbation of edge magnetic field

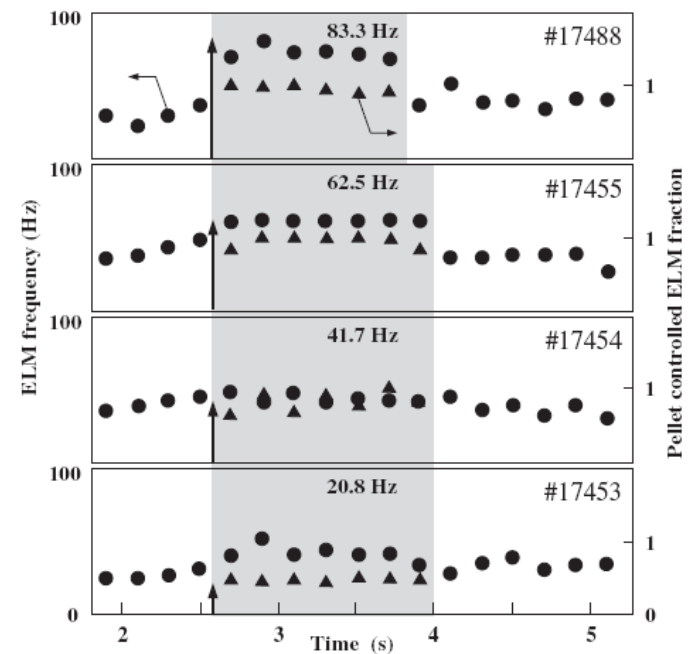
➤ Injection of frozen pellets → frozen hydrogen vaporizes leading to local  $n_e$  increase & (after thermalisation) pressure increase

➤ Local pellet perturbation → MHD instability & ELM

→  $f_{\text{pellet}} = f_{\text{ELM}}$  &  $\Delta W_{\text{ELM}}$  decreases in controlled way



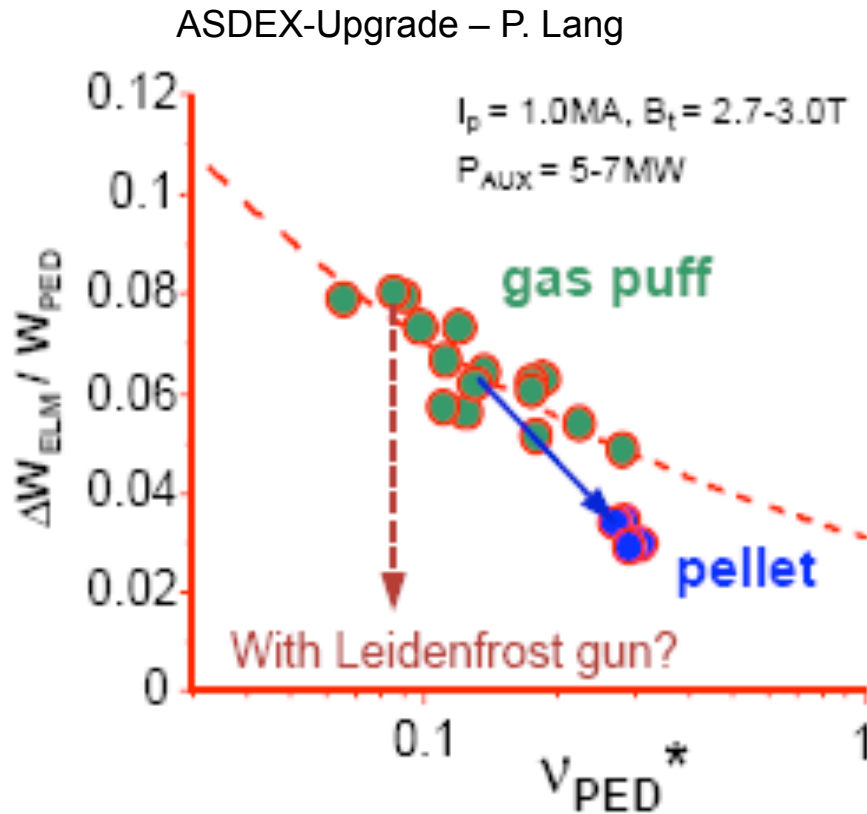
ASDEX-Upgrade – P. Lang NF'04



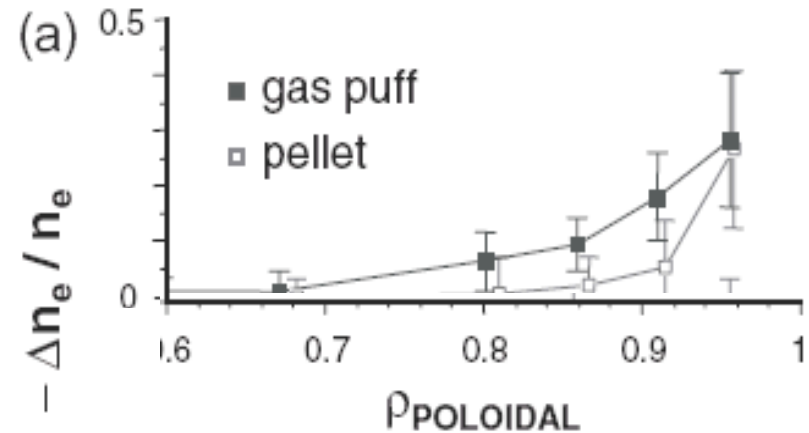
# Active control of ELMs (II)

Injection of frozen D pellets allows control of ELM trigger and  $\Delta W_{ELM}$  (ASDEX Upgrade – P. Lang)

Main effect associated with reduction of particle losses for pellet triggered ELMs



ASDEX Upgrade - Urano PPCF 2004



Present results far from ITER requirements (by factor of  $\sim 10$ )  $\rightarrow$  R&D with specially designed pellet injection systems/ experiments on going

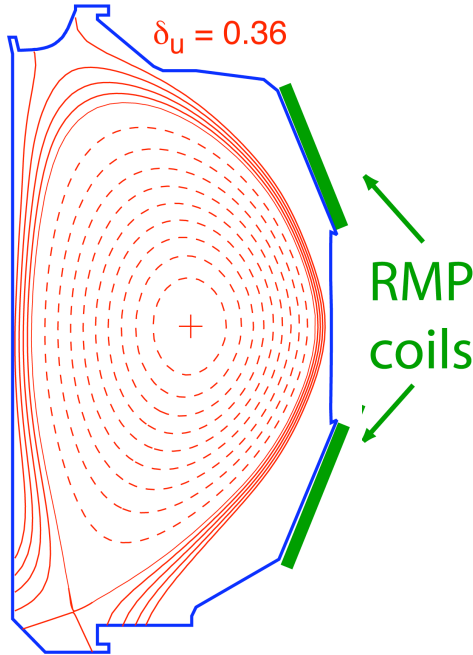
## Active control of ELMs (III)

- Applying an external magnetic perturbation of edge  $B_\theta$  allows complete control and/or suppression of ELMs
- Edge field lines are ergodised and edge plasma energy transport is affected

DIII-D- T. Evans NF'08

126006 3500 ms

$\delta_u = 0.36$

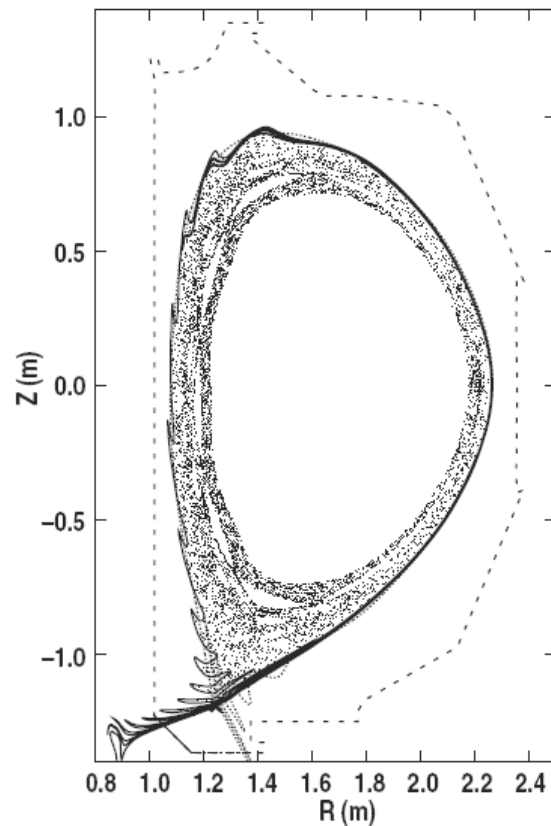


$\delta_{low} = 0.70$

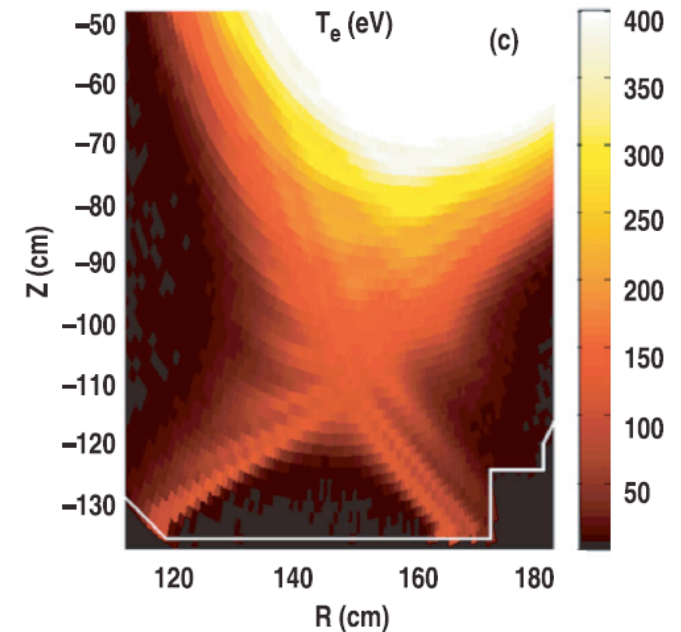
2006 ITER

Similar Shape

Wingen-PoP 2009

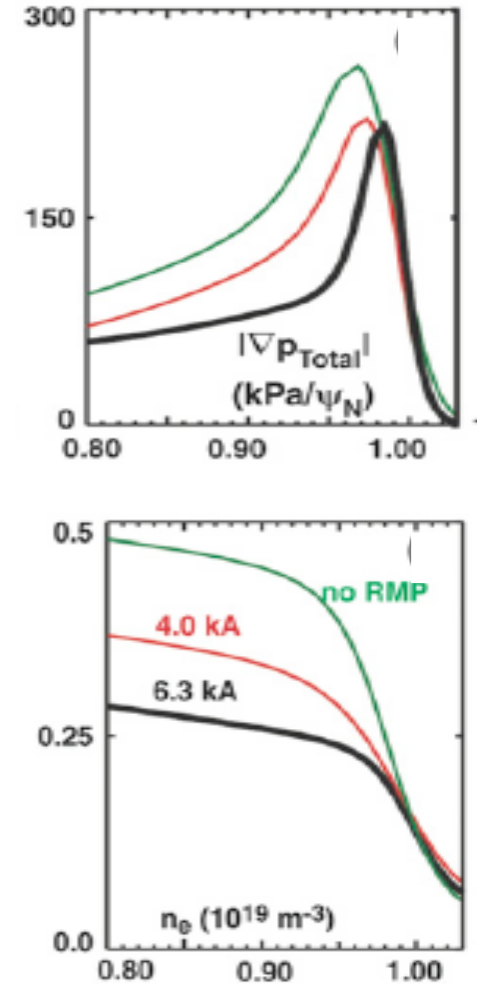
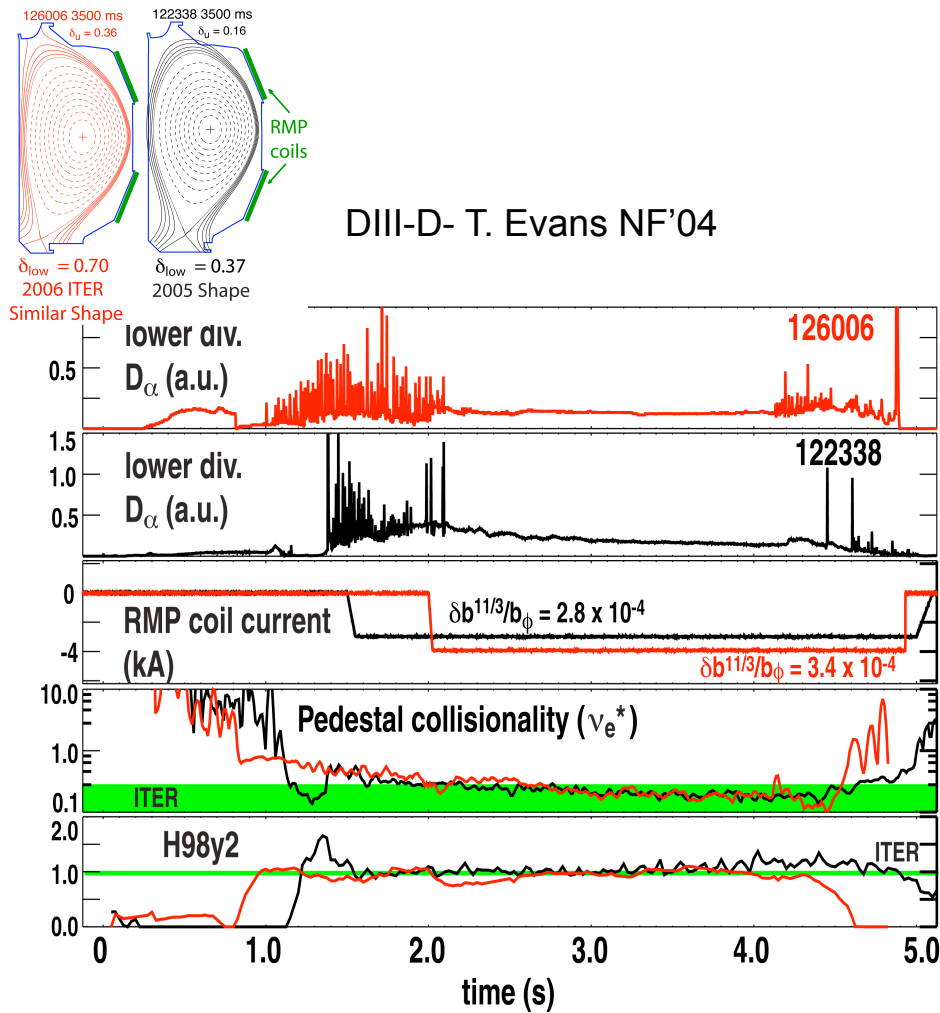


Joseph-NF 2008



# Active control of ELMs (IV)

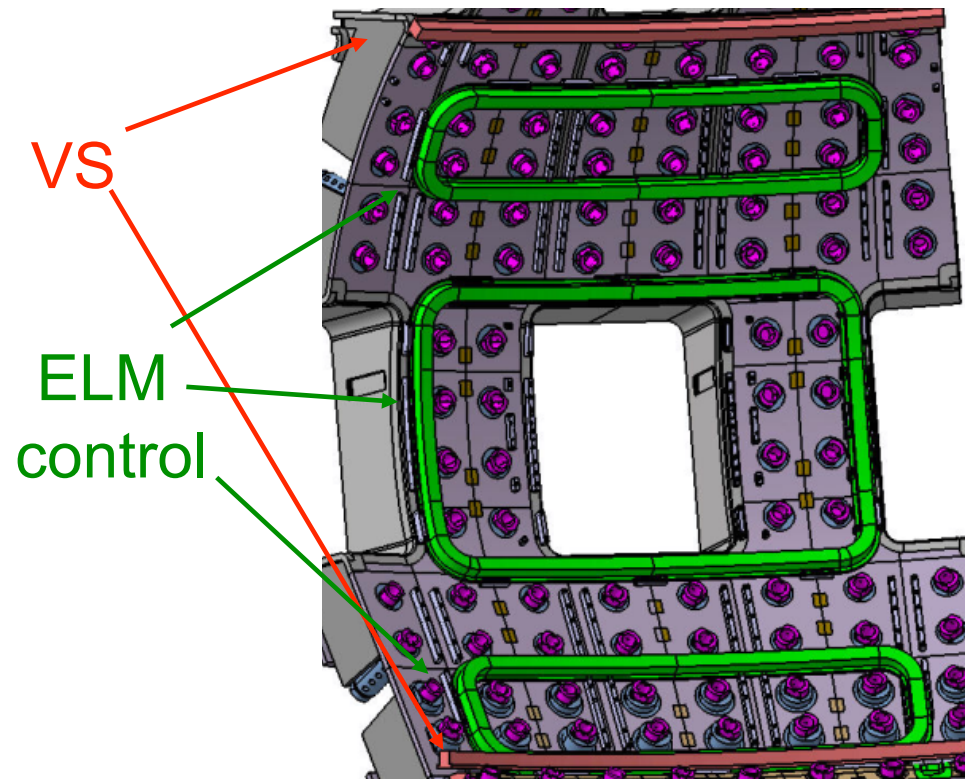
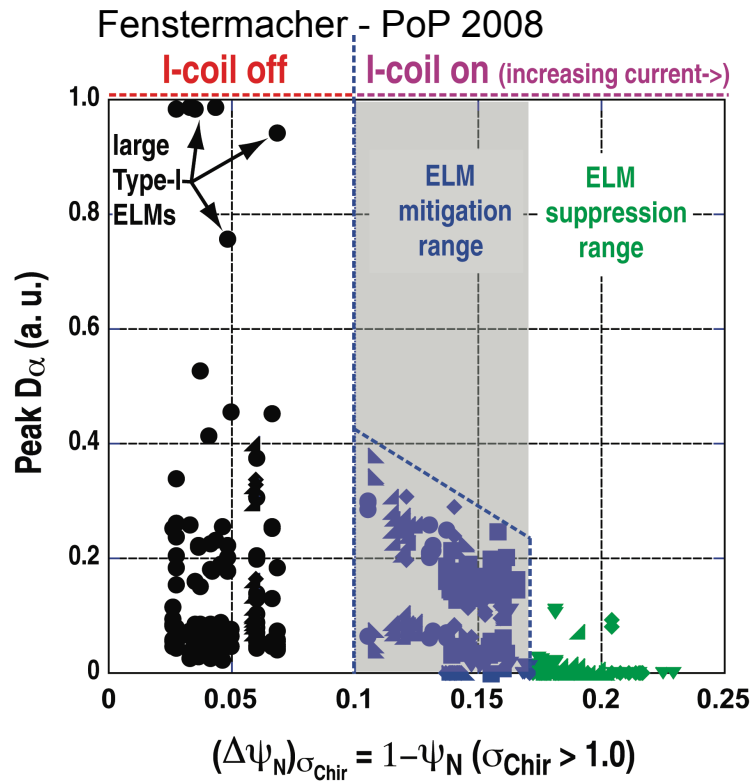
- With large enough  $B_\theta$  perturbation  $\rightarrow$  edge grad-p decreases & ELMs suppressed
- Major effect is  $n_e$  reduction not  $T_e \rightarrow$  experimental and theoretical R&D





# Active control of ELMs (V)

- ELM control coils for ITER designed on DIII-D based criterion for ELM suppression (~ 90-100 kAt, including 20% margin)
- Flexible system → all coils powered independently
- Technology R&D work ongoing to define coil conductor/insulation design
- Integration with other ITER components being finalised to fix interfaces



# Phenomena causing transients in ITER (I)

➤ Largest energy transients in ITER → disruptions

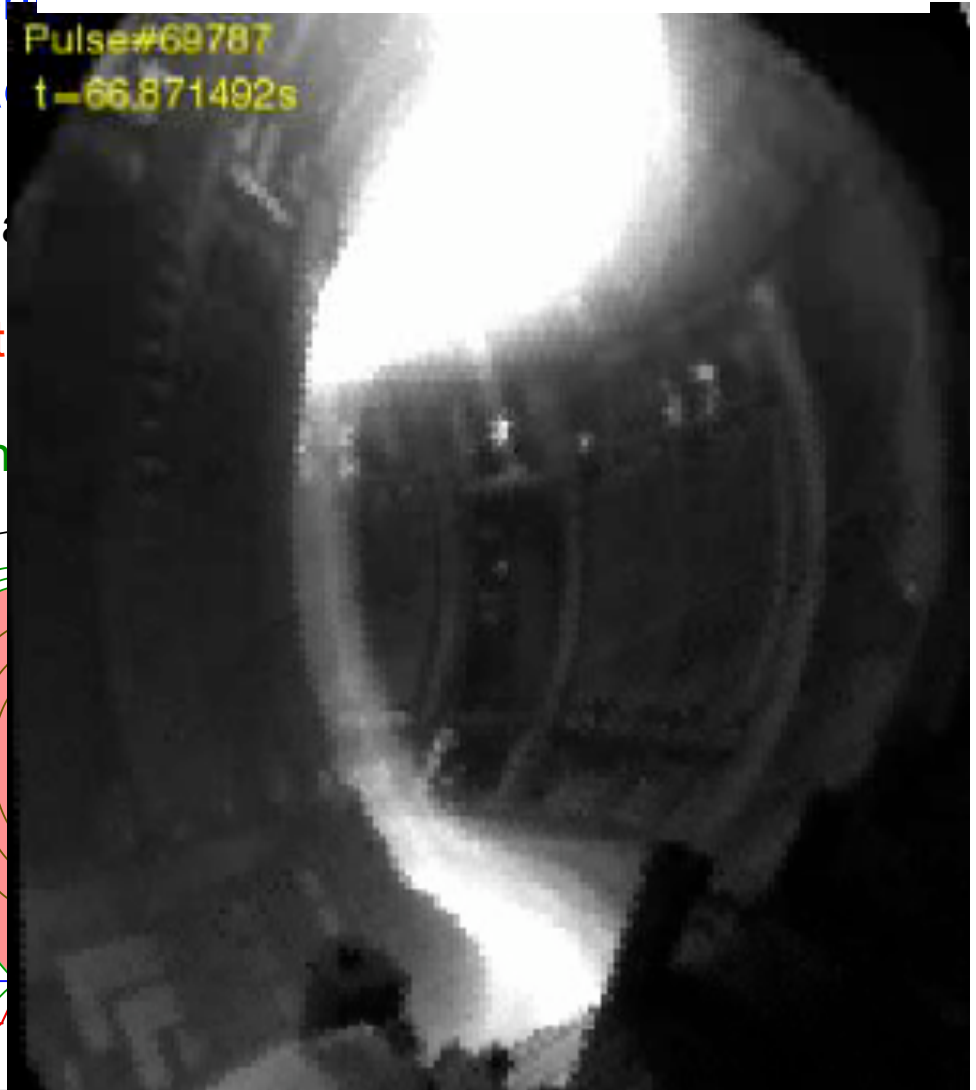
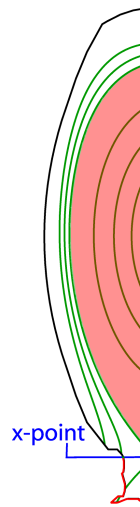
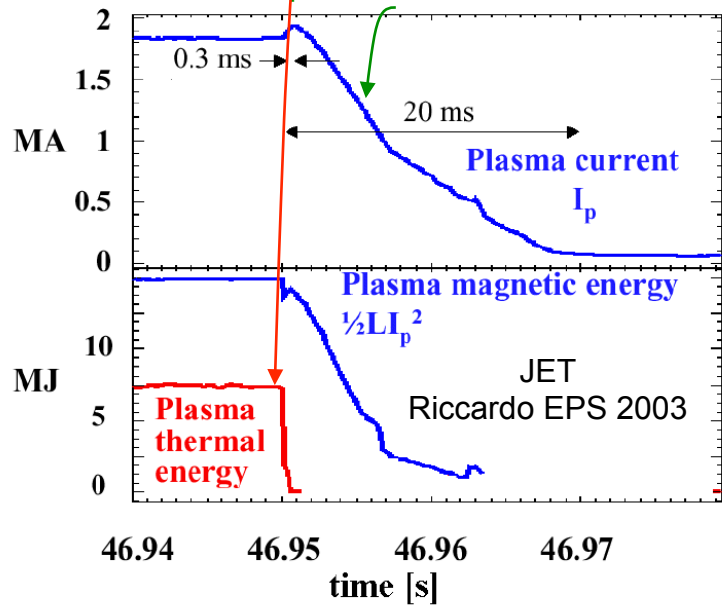
- ✓  $W_{\text{plasma}}$  → deposited by plasma on
- ✓  $W_{\text{magnetic}}$  → conductors & VV + rad

JET- J.A. Alonso

Plasma develops unstable  $p(r)$ ,  $j(r)$  → L

plasma confinement is destroyed (t

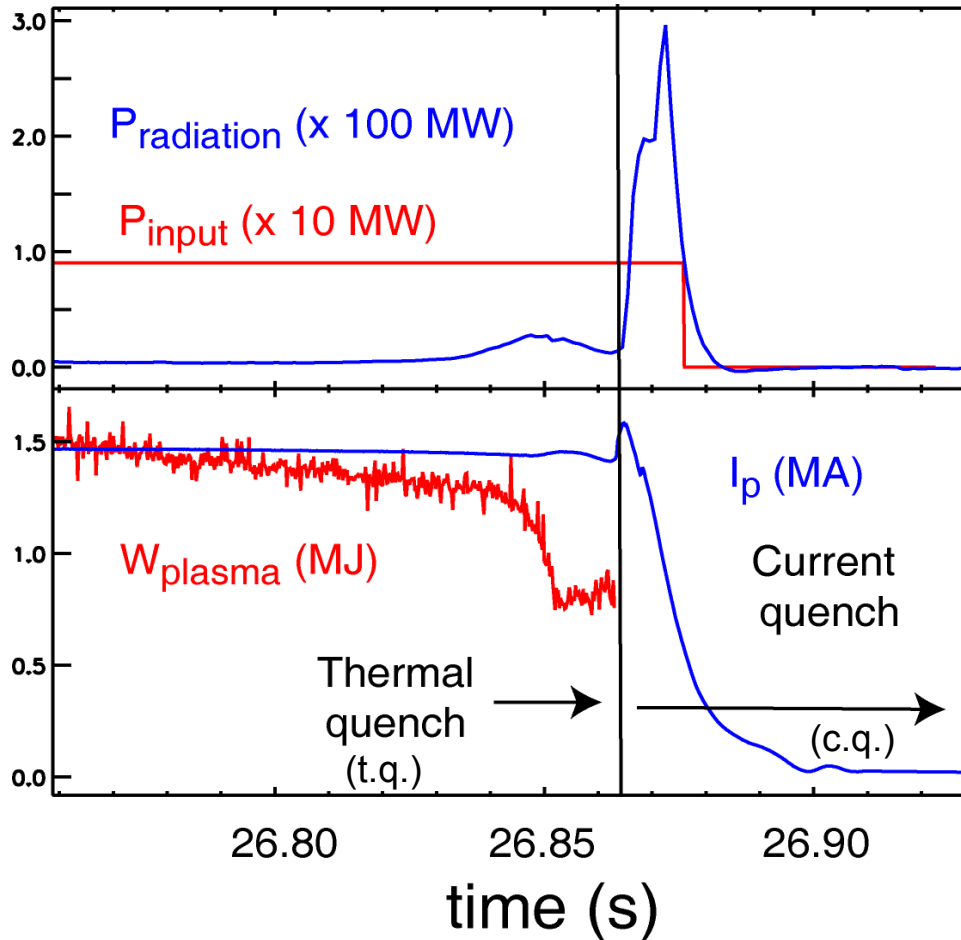
plasma current vanishes (current



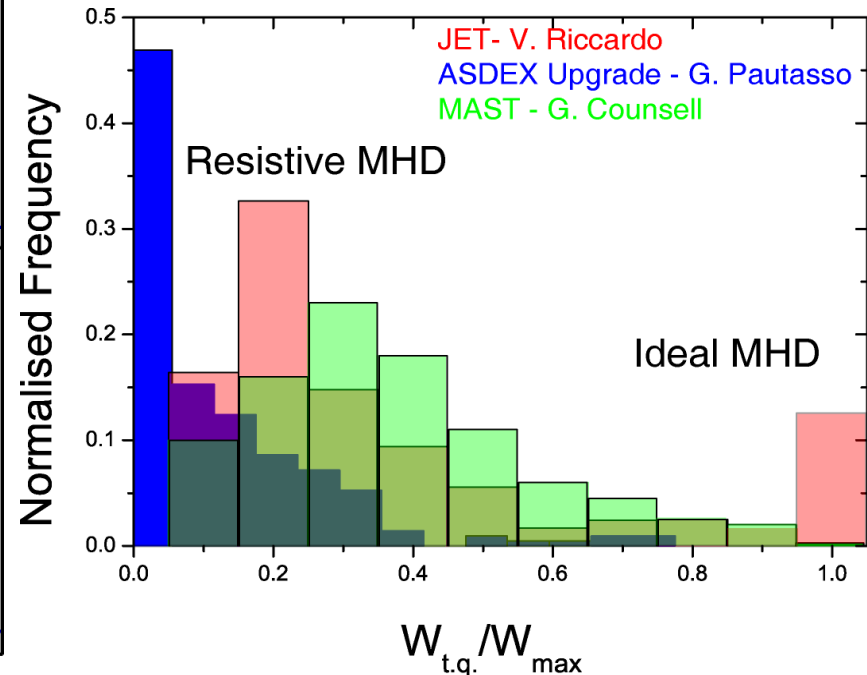
# Power fluxes to PFCs during disruptions (I)

- Thermal quench leads to largest disruptive power fluxes because of  $W_{\text{plasma}}$  and short timescale (excluding runaway electrons)
- $\tau_E$  deteriorates in advance of disruption  $\rightarrow W_{\text{t.q.}} \sim 0.3 W_{\text{plasma}}^{\text{H-mode}}$

JET-Pulse No. 69787

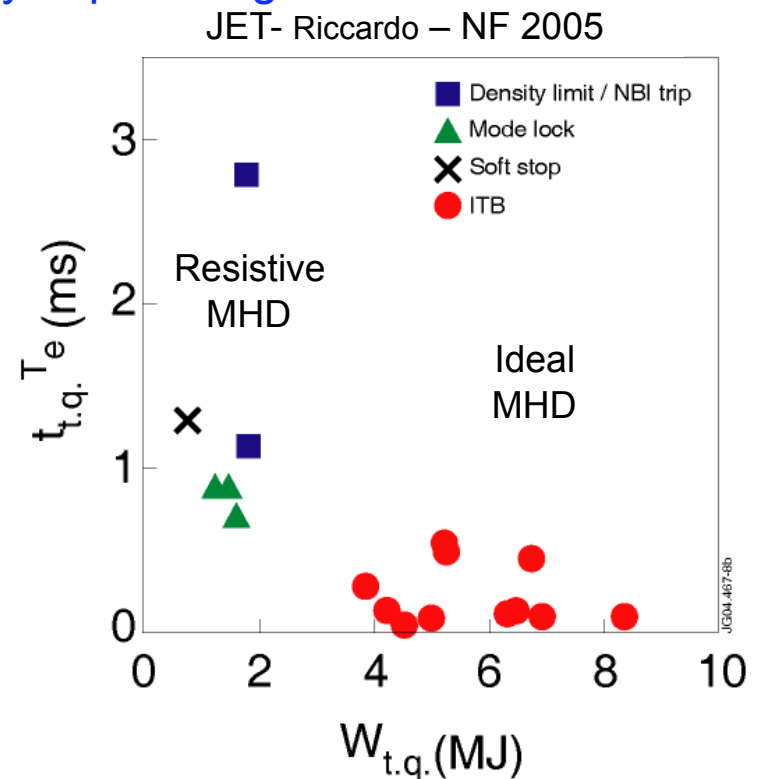
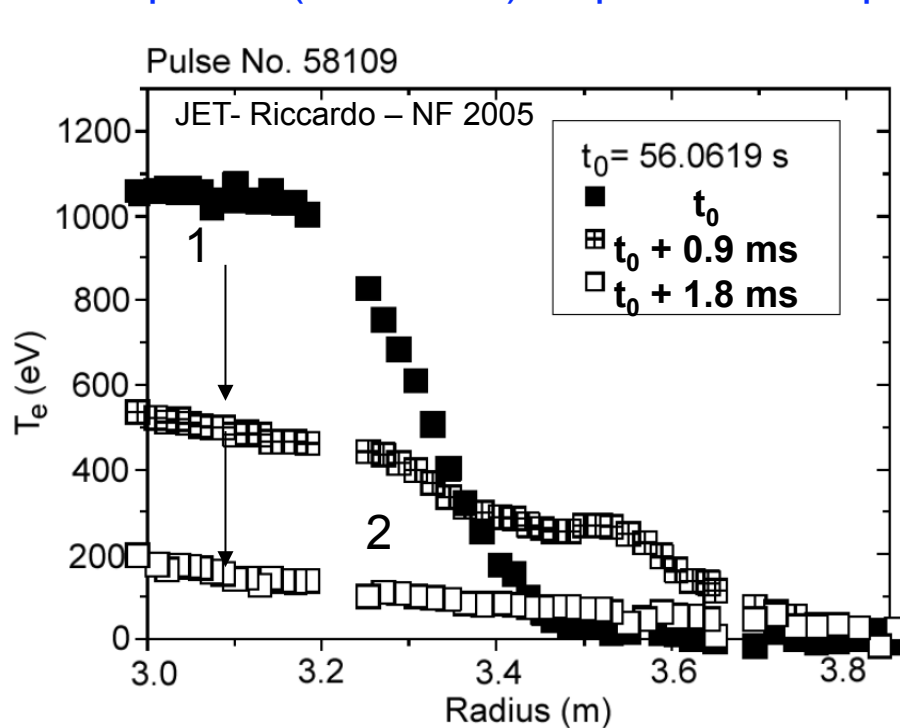


A. Loarte - EPS 2007



# Power fluxes to PFCs during disruptions (II)

- Power fluxes during thermal quench show large variability → complex processes leading to final plasma thermal energy collapse
- A. Disruptions (resistive MHD) :
  1. Enhanced transport in plasma core
  2. Loss of remaining plasma energy → flattening of current profile
- B. Disruptions (ideal limit) → plasma collapse by explosive growth of modes



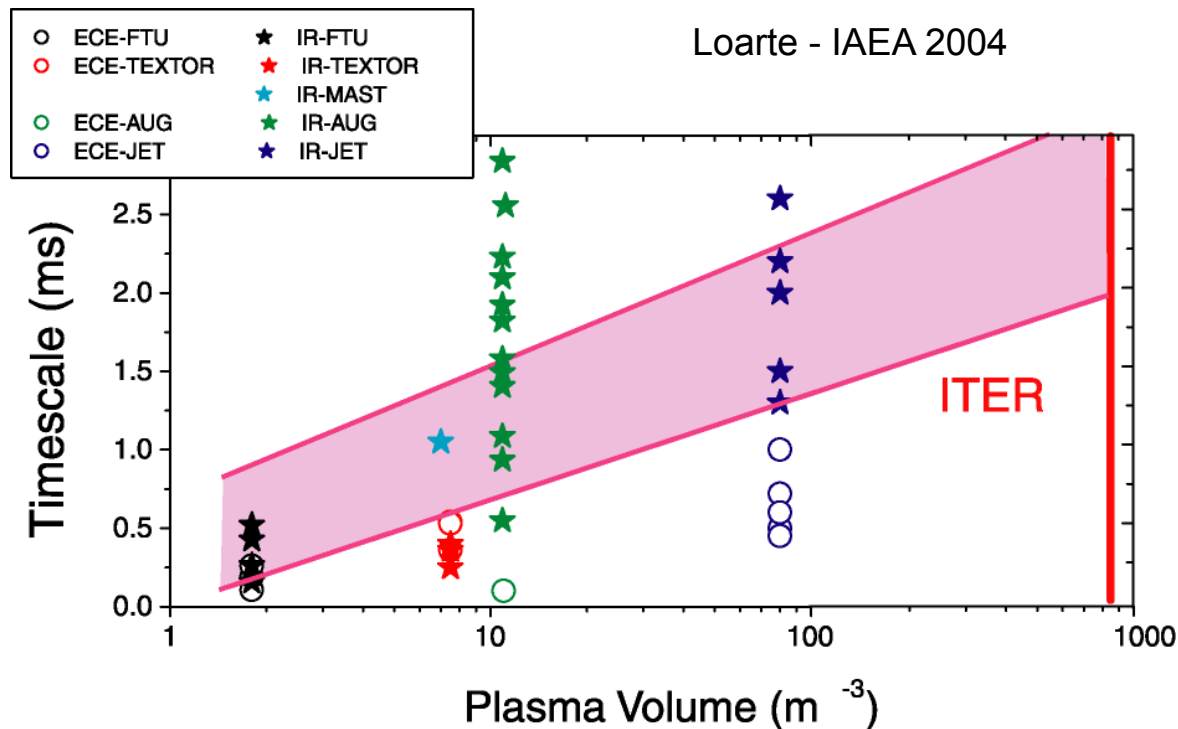
Resistive MHD plasma collapse → Longer timescales than ideal MHD collapse

## Power fluxes to PFCs during disruptions (III)

➤ Contrary to ELMs power flux timescales gets longer with size of device not inconsistent with energy diffusion in a strongly perturbed field

✓  $t_{t.q.} \sim R^2/\chi_{R-R}$

✓  $\chi_{R-R} \sim v_e L_c (\delta B_r/B)^2$  with  $L_c \sim R$  &  $v_e \sim T_e^{1/2} \sim R^\alpha$  ( $\alpha < 1$ )

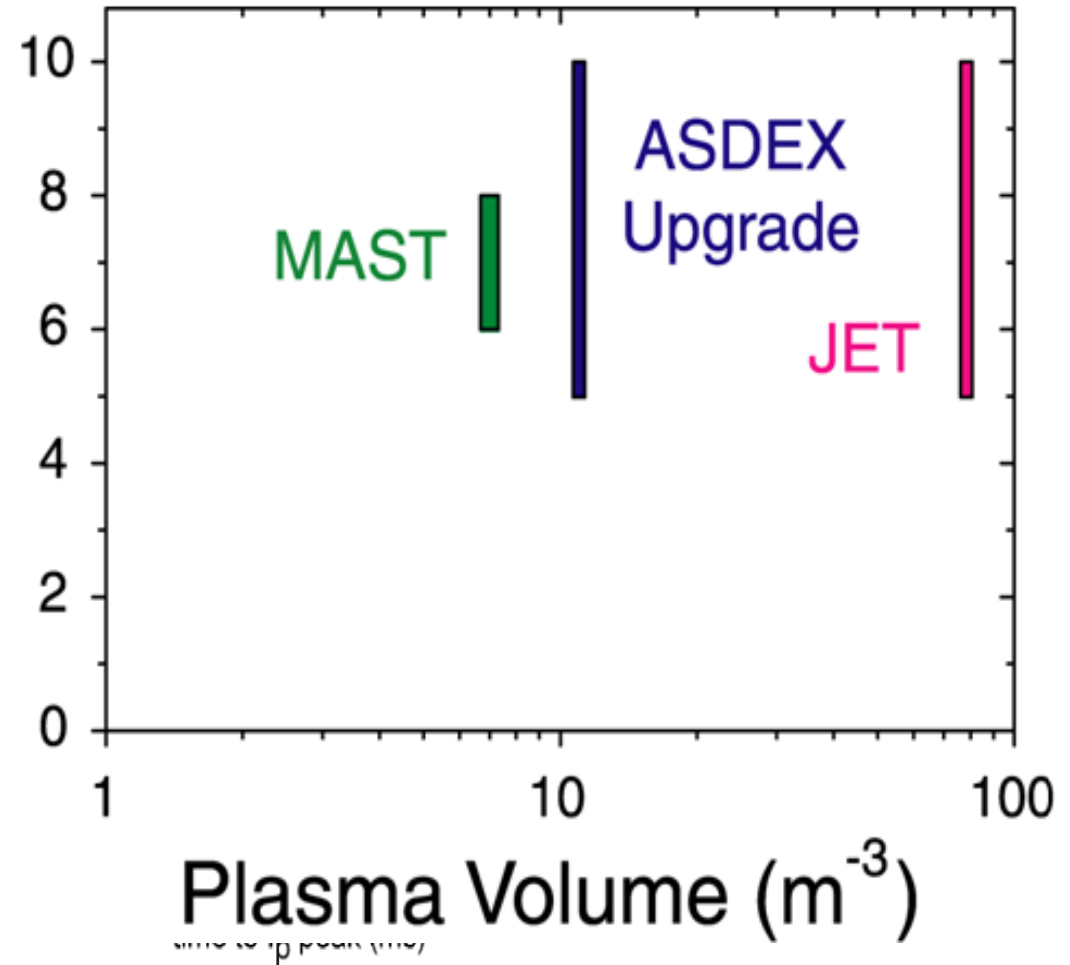
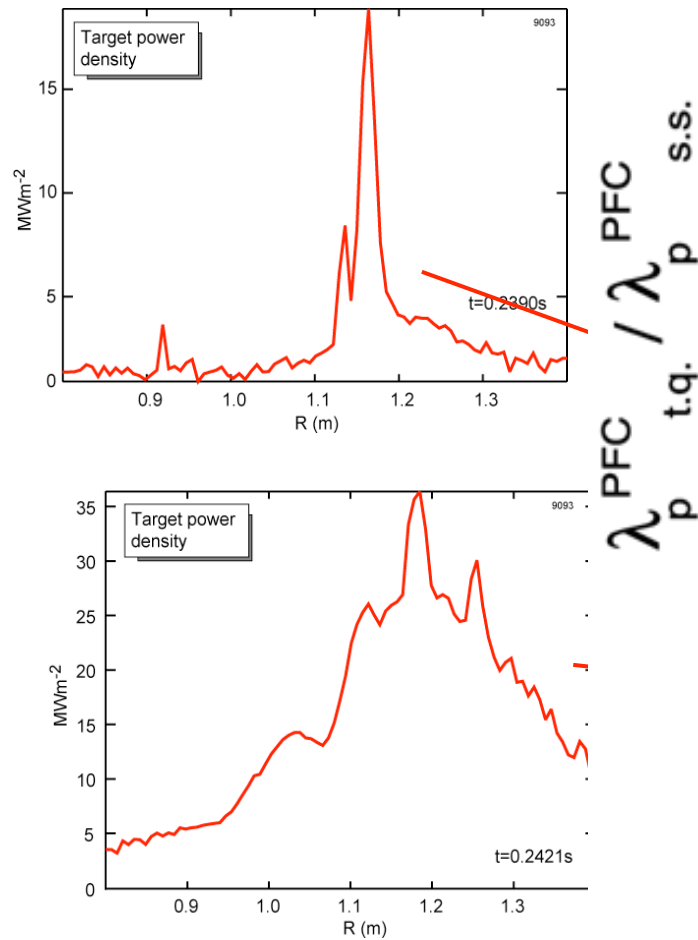


But large variability within a single experiment due to complex plasma dynamics during disruptions → no correlation with  $n_e, T_e$  before thermal quench

# Power fluxes to PFCs during disruptions (IV)

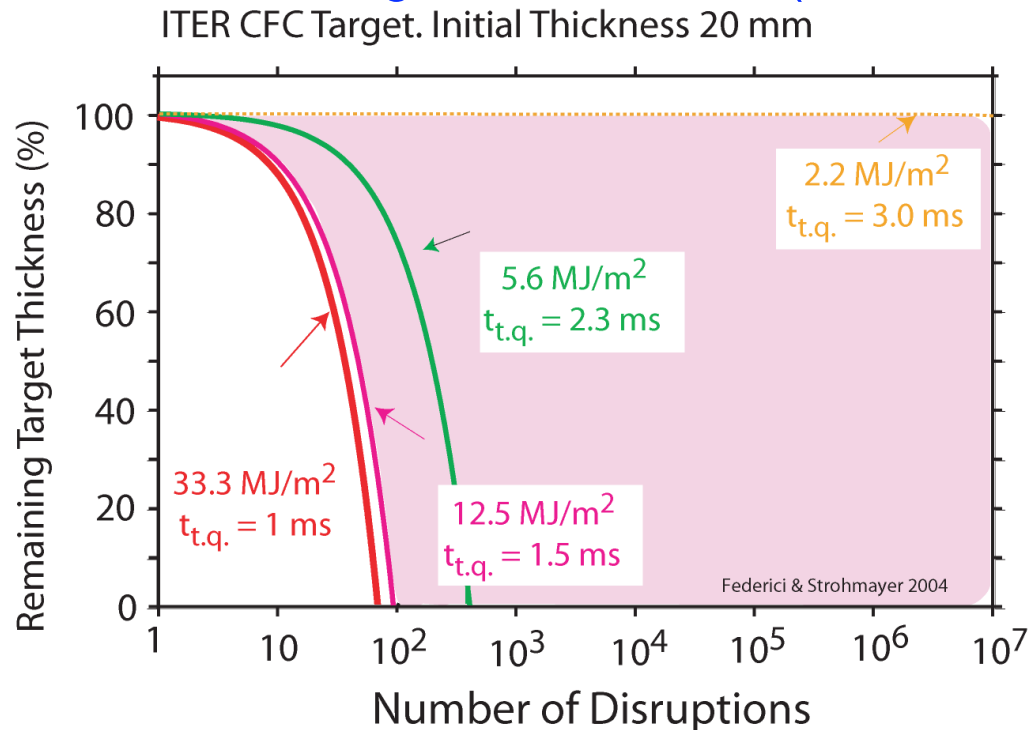
- Large broadening of the power flux footprint on PFCs (divertor targets) at thermal quench in divertor tokamaks → large  $\perp B$  transport (ergodisation of flux surfaces)

Loarte - IAEA 2004



## Power fluxes to PFCs during disruptions (V)

- Physics model and physics based extrapolations applied to determine disruption power fluxes to divertor and wall for thermal quench in ITER (R&D on-going)
- Expected maximum values similar to uncontrolled ELMs ( $\sim 20 \text{ MJm}^{-2}$ ) but over much large area of PFCs (5 -10 with respect to ELMs)



100 disruptions  $\rightarrow \sim 100 \mu\text{m}/\text{disruption}$   
 $1\text{m}^2$  erosion  $\rightarrow 200 \text{ g of C}$   
 ITER DT Plasma mass = 0.4 g

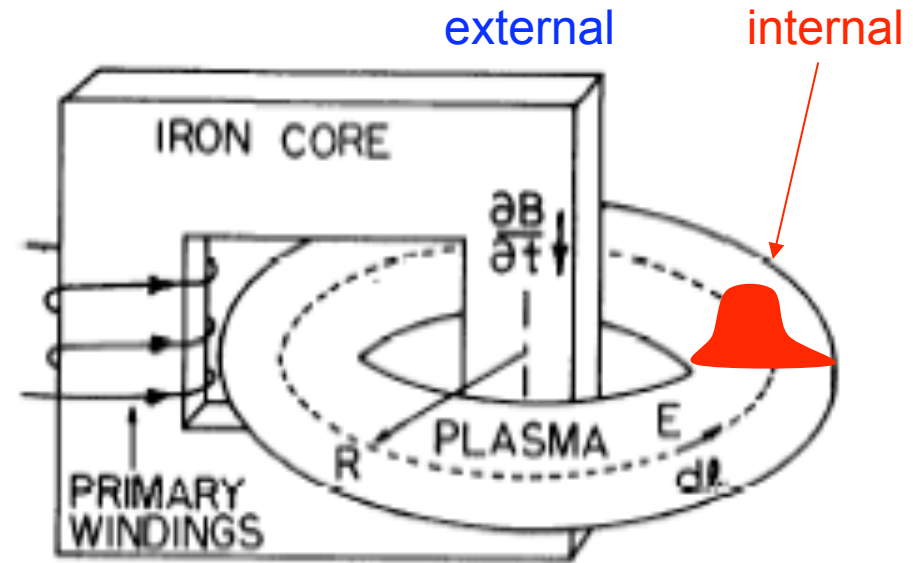
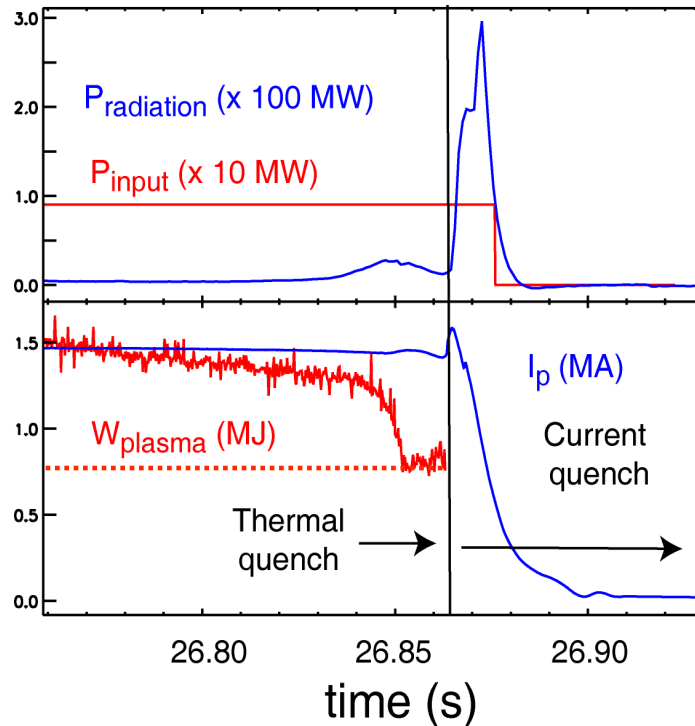
- Divertor lifetime  $\sim 100\text{s}$  high energy disruptions  $\rightarrow$  control of power fluxes during disruptions required for high performance ITER operation

# Power fluxes to PFCs during disruptions (VI)

- At thermal quench  $T_{\text{plasma}} \sim 10\text{s eV} \rightarrow$  plasma becomes resistive  $\eta \sim T^{-3/2}$
- V transformer small  $\rightarrow I_p$  decays (current quench)
- Internal magnetic energy  $\rightarrow$  Joule heating  $\rightarrow$  radiation by partly ionised impurity ions from thermal quench

JET-Pulse No. 69787

$$W_{\text{magnetic}} = \frac{1}{2} L_{\text{plasma}} I_p^2 \quad \text{with} \quad L_{\text{plasma}} = \underbrace{\mu_0 R_0 \left( \ln \left( 8 \frac{R_0}{a} \right) - 2 \right)}_{\text{external}} + \underbrace{\frac{\ell_i}{2}}_{\text{internal}}$$

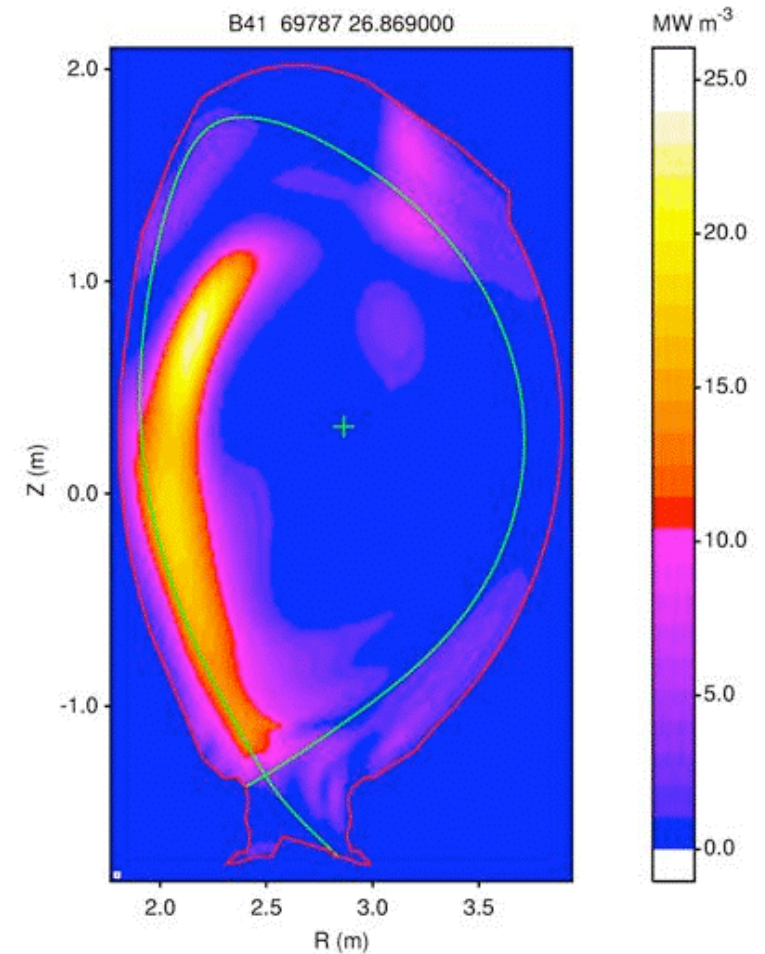
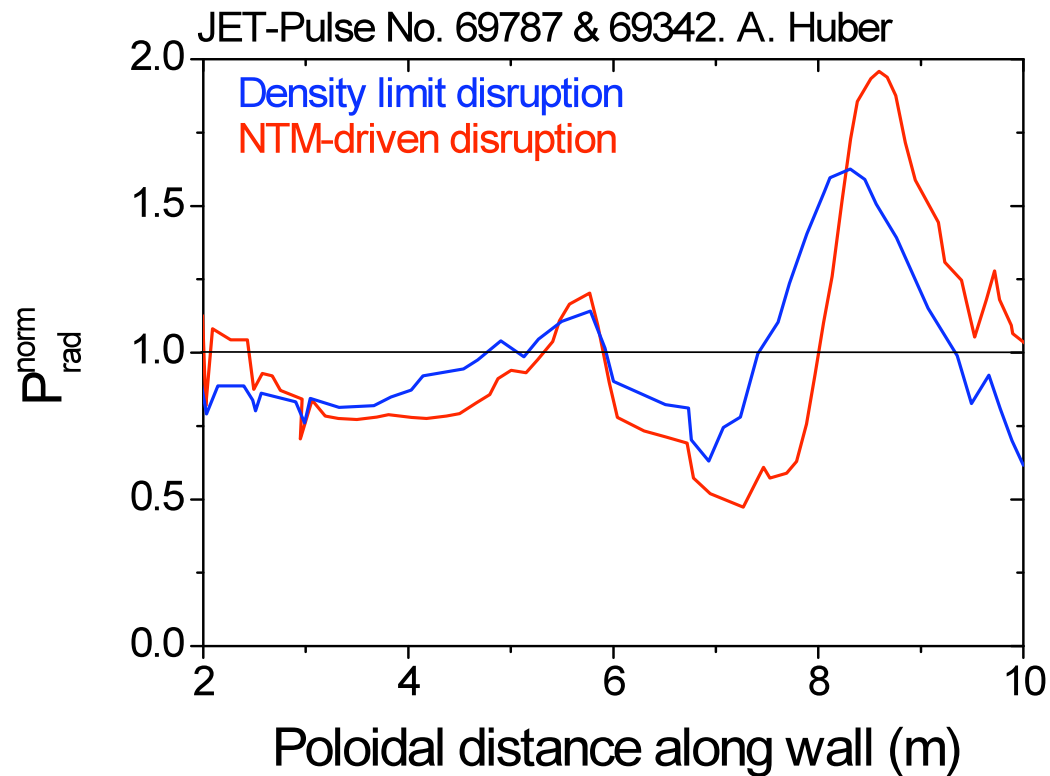


- Most external magnetic energy coupled back to external conductors by induction



# Power fluxes to PFCs during disruptions (VII)

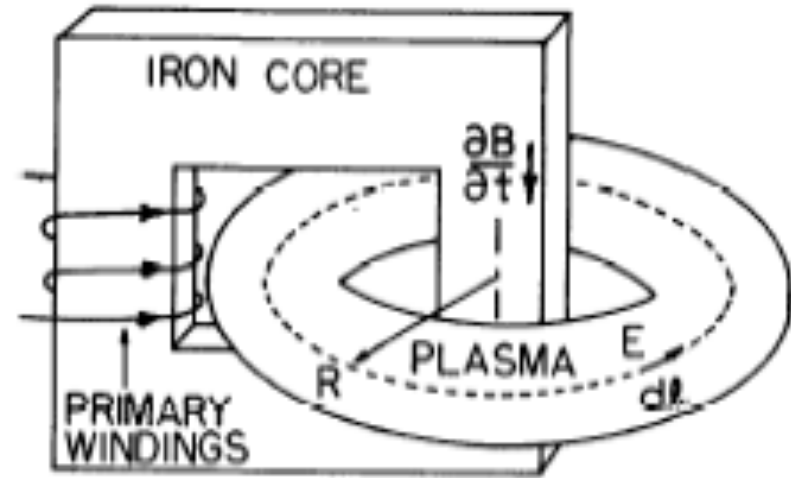
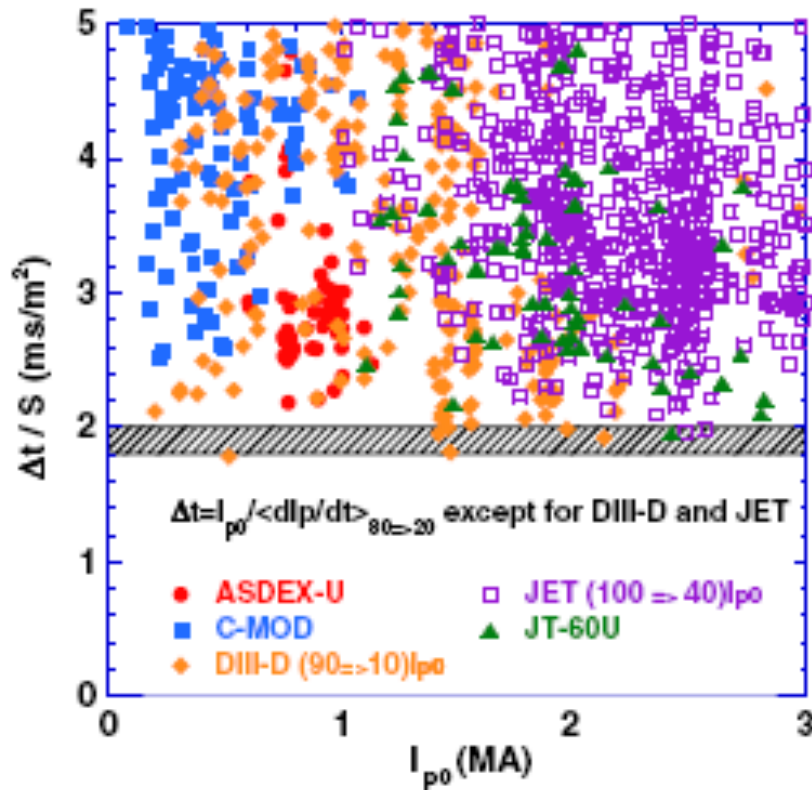
- $W_{\text{ohmic}} = W_{\text{mag}} - W_{\text{conductors}} \rightarrow$  plasma heating and radiation
- Radiation to distributed power flux (from magnetic energy loss) during current quench



# Power fluxes to PFCs during disruptions (VIII)

- Timescale of current quench scales with device size → 20 – 40 ms in ITER
- Long time scale and distribution of power by radiation → relatively low energy flux on PFCs during current quench
- Large induced E → runaway electrons

PIPB – NF 2007



$$\Phi_{\text{magnetic}} = L_{\text{plasma}} I_p$$

$$\frac{d\Phi_{\text{magnetic}}}{dt} = L_{\text{plasma}} \frac{dI_p}{dt}$$

$$\oint_{\partial S} \mathbf{E} \cdot d\mathbf{l} = -\frac{\partial \Phi_{B,S}}{\partial t}$$

## Power fluxes to PFCs during disruptions (IX)

- Electric field induced in plasma during current quench

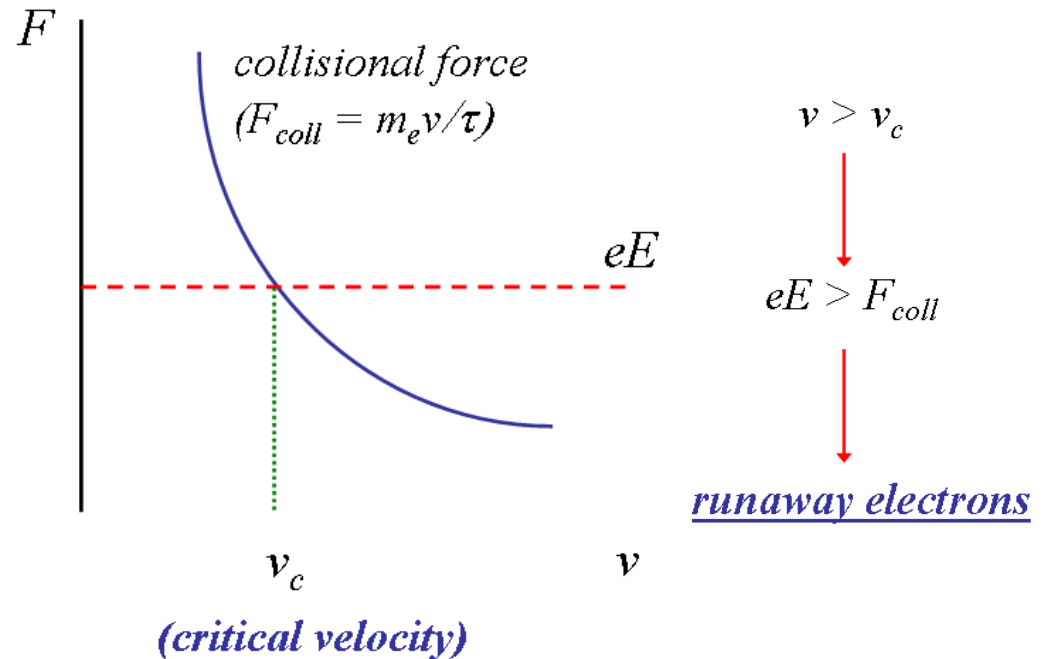
$$E_{\parallel} 2\pi R_0 = \frac{1}{2} \mu_0 R_0 \ell_i \frac{dI_p}{dt} \sim \frac{1}{2} \mu_0 R_0 \ell_i \frac{I_p}{CS_{plasma}}$$

- If field large enough → some electrons accelerated to  $v_e \sim c$  (runaway  $e^-$ )  
 $e^-$  in plasmas subject to acceleration (by  $E$ ) and deceleration by collisions

$$F_a = eE$$

$$F_d = m_e \frac{v_e}{\tau_s} \rightarrow F_d = m_e \frac{v_e}{\tau_s} \propto \frac{1}{v_e^2}$$

$$\tau_s = \left( \frac{4\pi\epsilon_0^2 m_e^2}{(Z^2 + 2)ne^4 \ln \Lambda} \right) * v_e^3 \propto v_e^3$$



## Power fluxes to PFCs during disruptions (IX)

- Formation of sizeable runaway current limited by diffusion in velocity space  
Dreicer mechanism → critical electric field dependent only on plasma density

$$E > \frac{ne^3 \ln \Lambda}{4\pi\epsilon_0^2 m_e c^2}$$

Runaway electron generation in current quench  
(no runaways if n sufficiently high)

- Collisions between runaways and thermal electrons can also create secondary runaways (avalanche)

Generation rate

$$R_f = \frac{1}{\tau_R} \quad \tau_R = 3 \ln \Lambda \frac{m_e c}{eE}$$

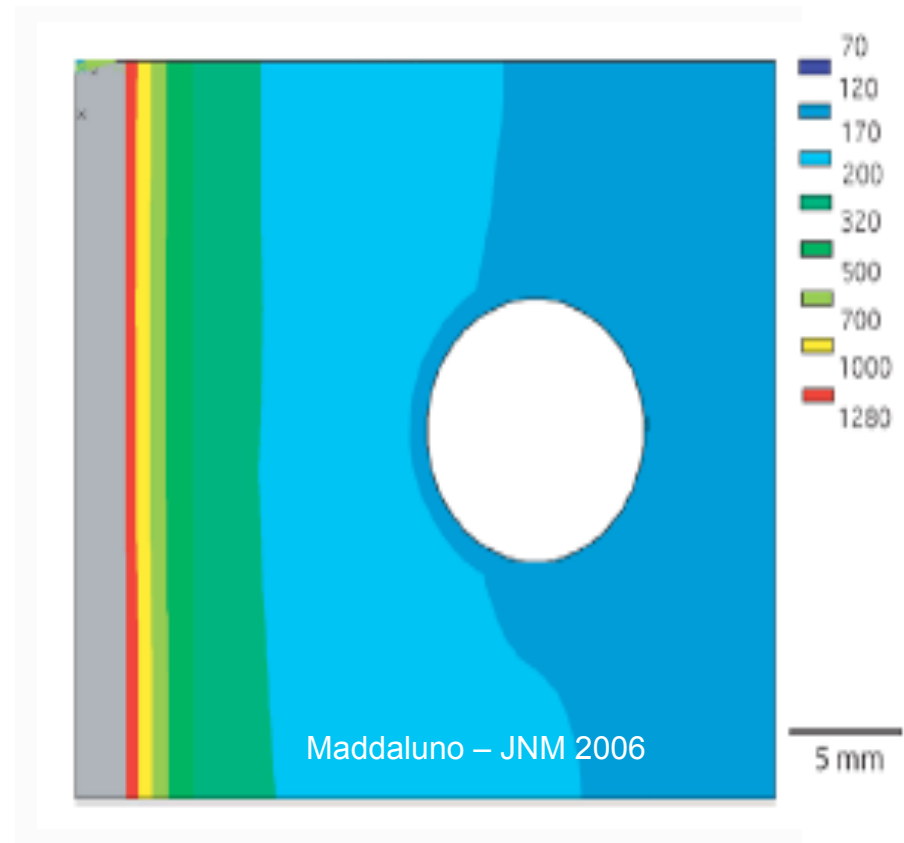
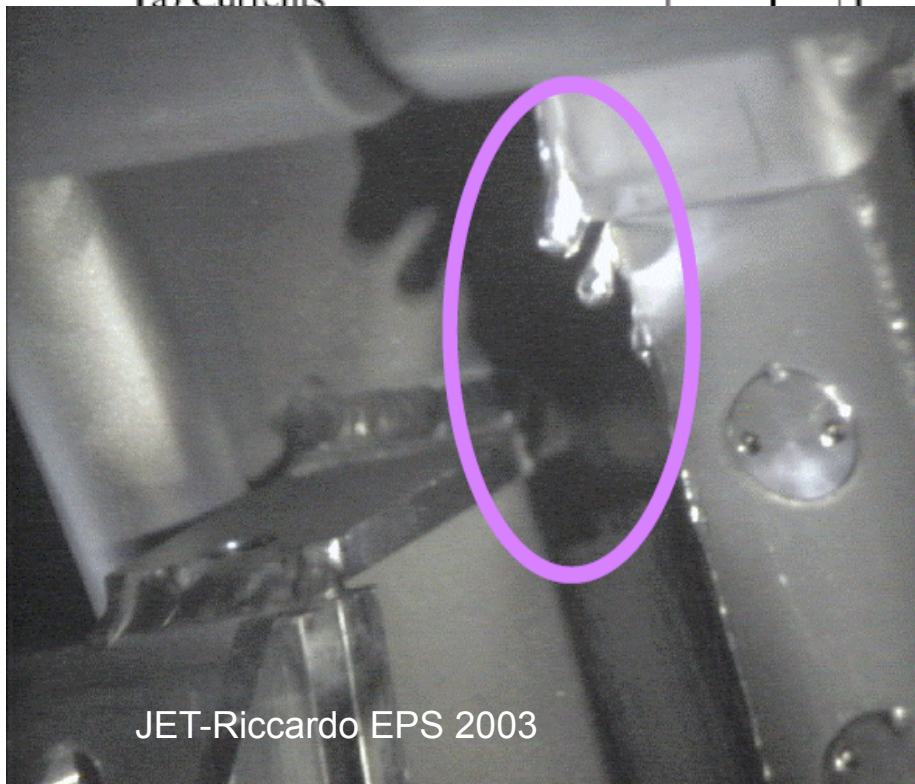
$$I_r = I_0 e^{t/\tau_r} \approx I_0 e^{2.5 I_p (MA)}$$

For ITER ( $I_p = 15$  MA)  $I_r \sim I_p$  before thermal quench with  $E_e \sim 10$ s MeV

## Power fluxes to PFCs during disruptions (X)

- Runaway electron discharges become vertically unstable and deposit their energy on localised areas of the first wall
- Deep melting (~ mm) expected in ITER (seen already in present generation of tokamaks) → problems for water cooled components (< 1 cm thick PFM)

Runaway electron damage in JET



runaway formation avoidance and/or controlled energy deposition required for ITER

# Mitigation of disruptions (I)

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Disruption mitigation schemes aim at reducing consequences of disruptions in ITER

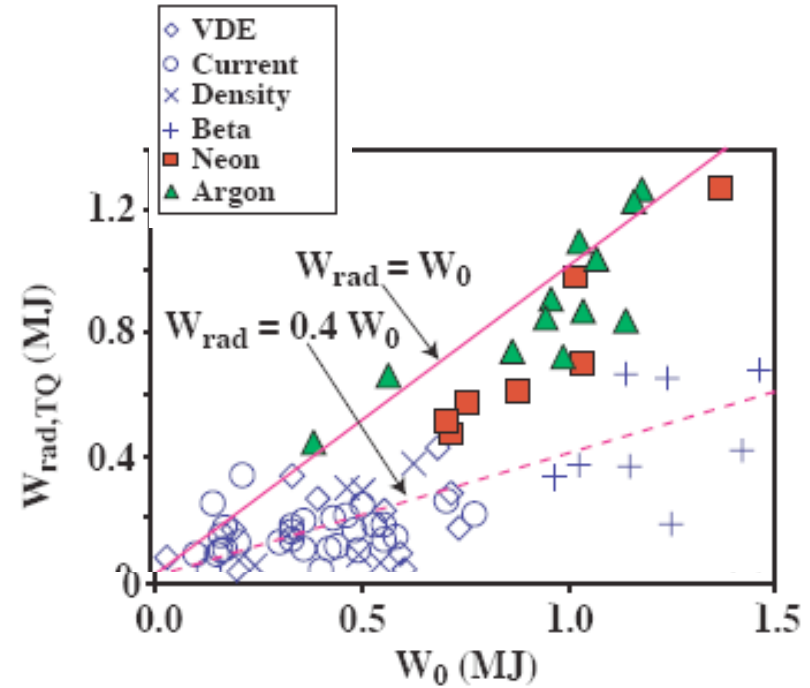
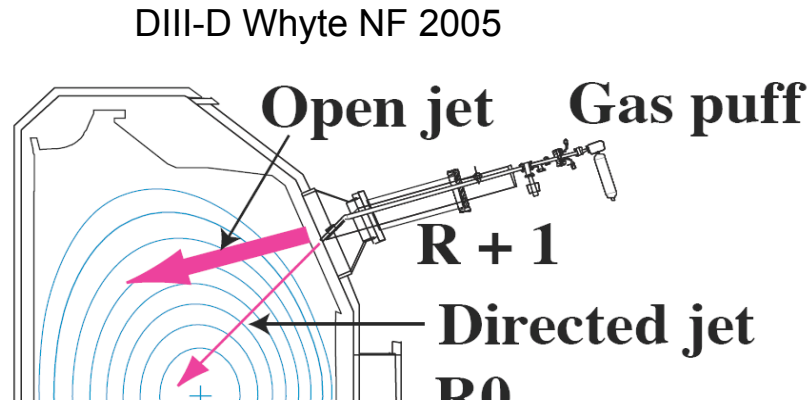
1. High power fluxes onto PFCs
2. Large forces on VV caused by plasma displacement before disruption
3. Formation and localised impact of runaways

Schemes foreseen rely on present experimental results

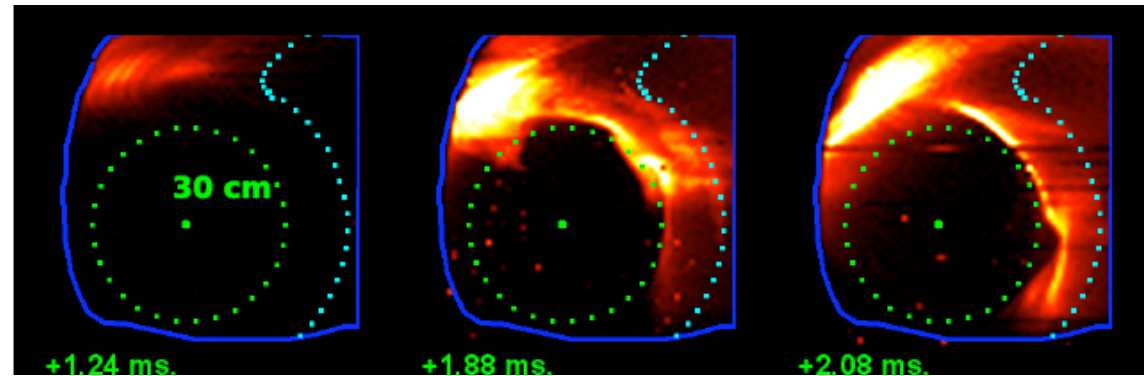
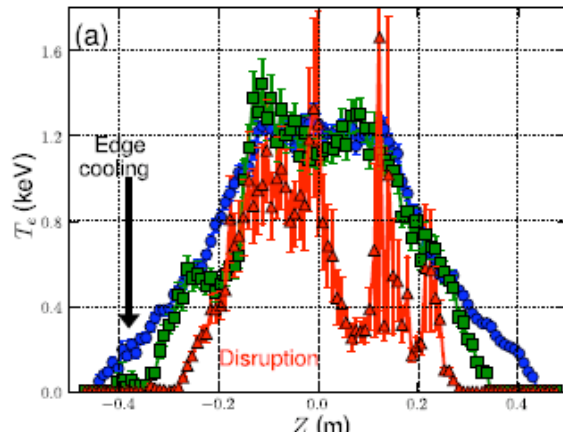
- Injection of large amount of material before disruption (OK for 1, 2 and possibly 3)
- Application of schemes for soft landing of runaway electron discharges (for 3)

# Mitigation of disruptions (II)

- Injection of large amount of impurities before disruption can radiate plasma energy at thermal quench → more distributed power load on PFCs

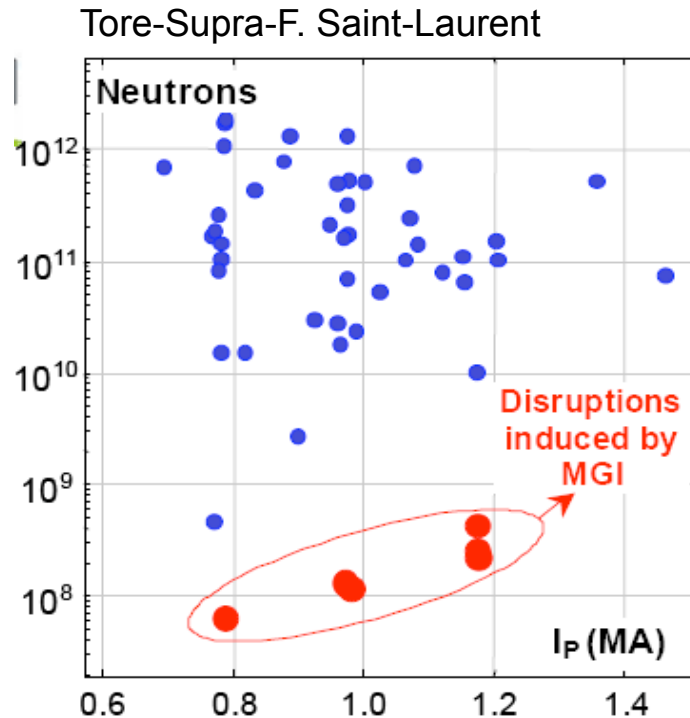


Bozhenkov TEXTOR PPCF 2008



## Mitigation of disruptions (III)

- Injection of large amount of gas is effective in also effective in suppressing runaway electrons formed during current quench in present experiments



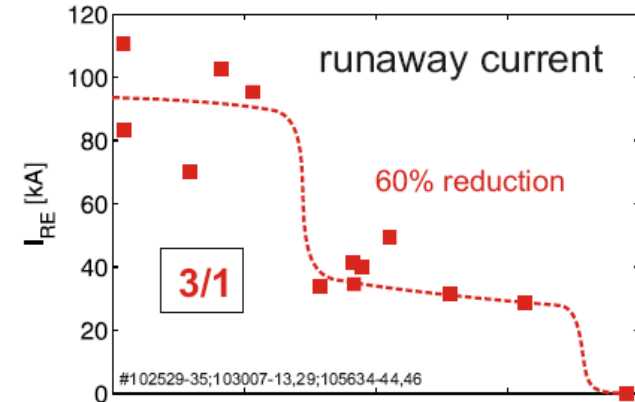
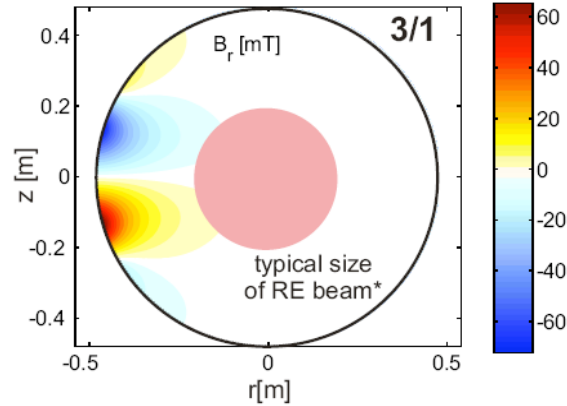
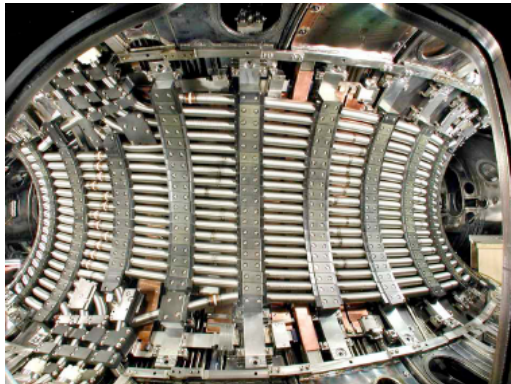
- Suppression of avalanche mechanism in ITER requires  $n_e \sim 5 \cdot 10^{22} \text{ m}^{-3} \rightarrow$  100's g of material injected in few ms
  - ✓ Complex technology
  - ✓ Complications in restoring high vacuum after mitigation



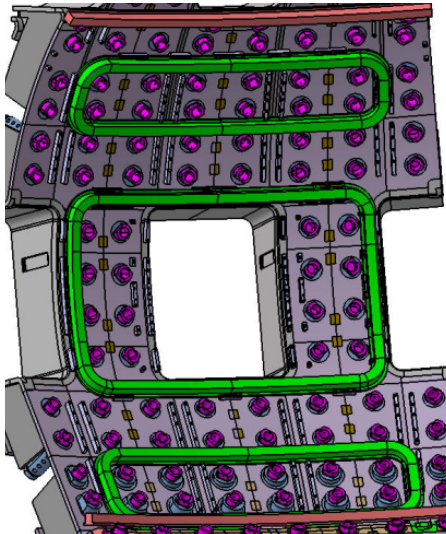
# Mitigation of disruptions (IV)

- Solution investigated → combination of massive material injection & enhanced runaway loss by perturbed magnetic field (Lehnen PRL'08)

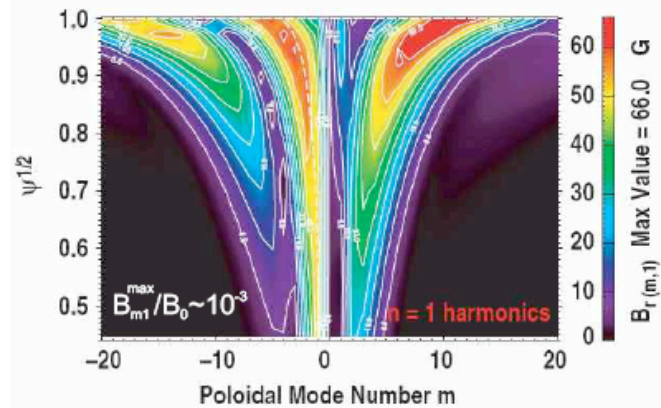
## TEXTOR



## ITER



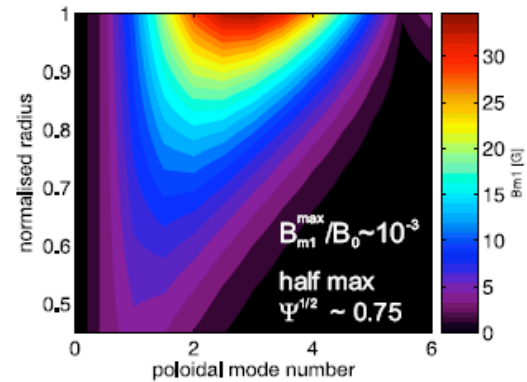
ITER study for n=1 mitigation



M.J. Schaffer et al, Nucl. Fusion 48 (2008) 024004

## TEXTOR n=1

$I_{DED} = 1.4\text{kA}$  (RE threshold)

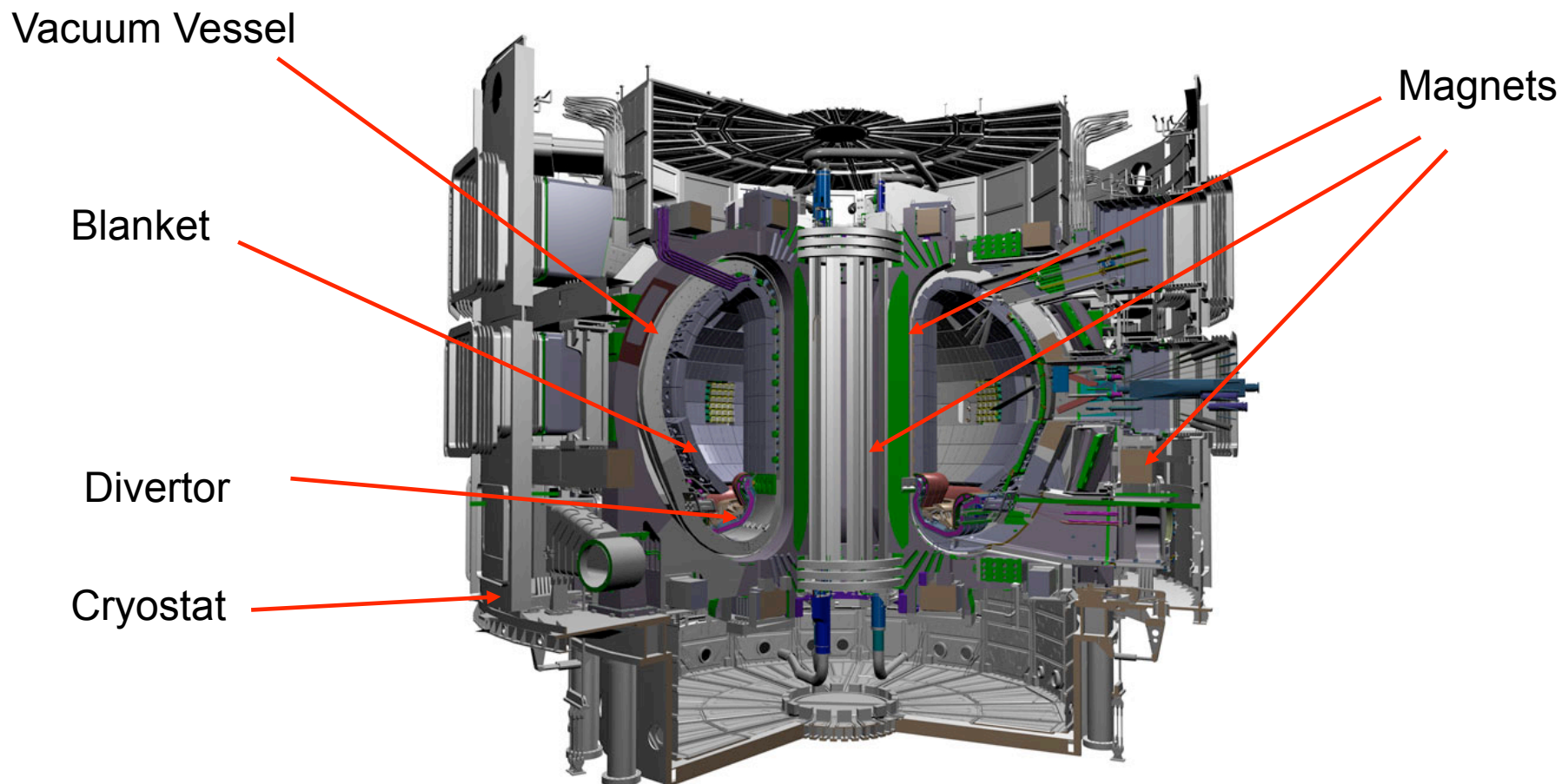


# Summary and Conclusions

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- Physics of power transients in ITER involves large range of physical processes not fully understood
  - Confined plasma MHD
  - Transport of energy and particles along distorted magnetic surfaces
  - Interaction of hot plasmas ( $\sim$  keV) with material surfaces
  - Formation of high energy electron plasmas ( $\sim$  MeV) and their interaction with material surfaces
- Understanding physics of power transients is required for their control and mitigation
- **Reliable ITER operation as required to achieve ITER's goals is synonymous of reliable control and mitigation of transients**

# Plasmas in ITER (II)



Diagnostics and H&CD systems (33 MW NNBI, 20 MW ICRH, 20 MW ECRH)

- Machine mass: 23350 t (Cryostat + VV + Magnets)
- Inductive operation:  $R = 6.2$  m,  $a = 2.0$  m,  $\kappa_{\text{sep}} = 1.85$ ,  $\delta_{\text{sep}} = 0.48$ ,  $I_p = 15$  MA,  $B_t = 5.3$  T

# Power fluxes to PFCs during ELMs (III)

- Physics of energy flow from plasma to PFCs → competition of parallel and perpendicular transport after linear MHD phase

